

# **Quantum Spintronics Design**

## **(NV centers in diamond)**

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2020.02.19

CMD Spintronics Design Course  
@Osaka University



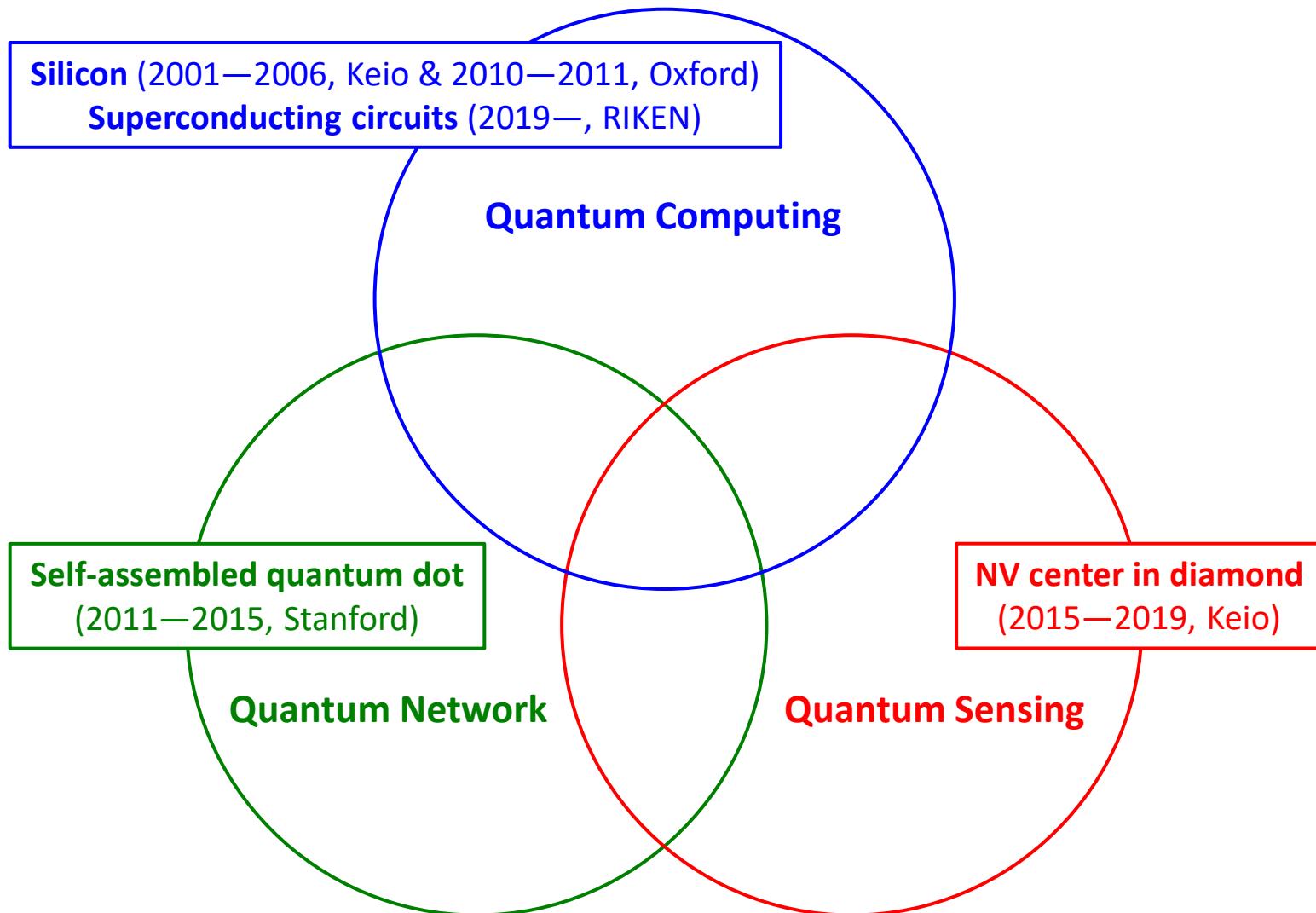
# Short CV



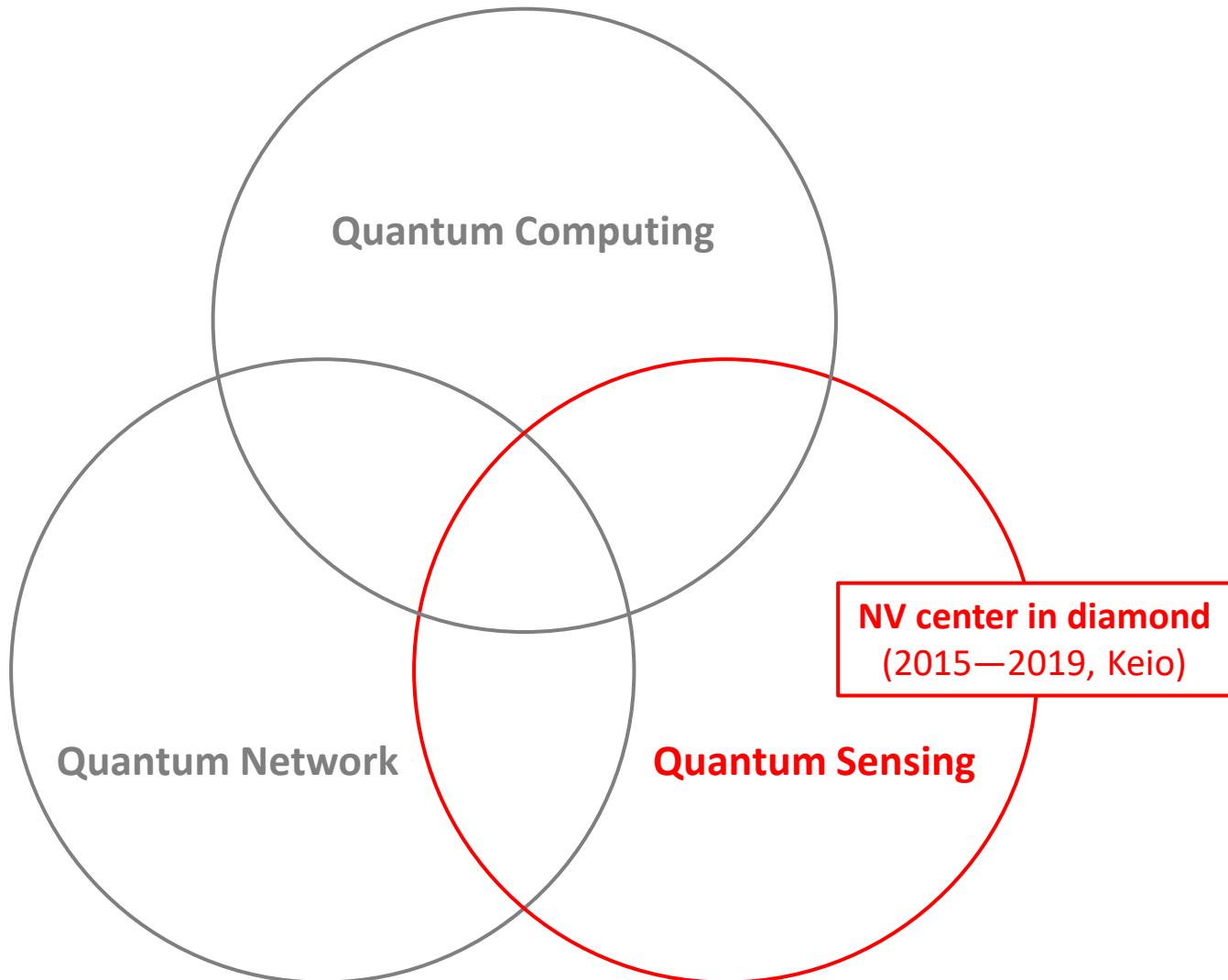
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- **2001.4 – 2006.3 (Keio)** → Quantum computing (silicon)
- **2006.4 – 2009.12 (ISSP, UT)** → Quantum transport (GaAs QDs, Josephson)
- **2010.1 – 2011.6 (Oxford)** → Hybrid system (spin–cavity coupling)
- **2011.7 – 2015.3 (Stanford/RIKEN)** → Quantum network (InAs QDs)
- **2015.4 – 2019.1 (Keio)** → Quantum sensing (diamond)
- **2019.2 – Present (RIKEN)** → Quantum computing (Josephson)

# Quantum technologies



# Quantum technologies



# Outline

- **Basics of NV centers in diamond**
  - Structure
  - Optical properties
  - Spin properties and control
- **Quantum sensing**
  - Principle of AC magnetometry
  - Detection of proton spin ensemble
  - Detection and localization of a single  $^{13}\text{C}$  nuclear spin
  - Ultrahigh resolution sensing

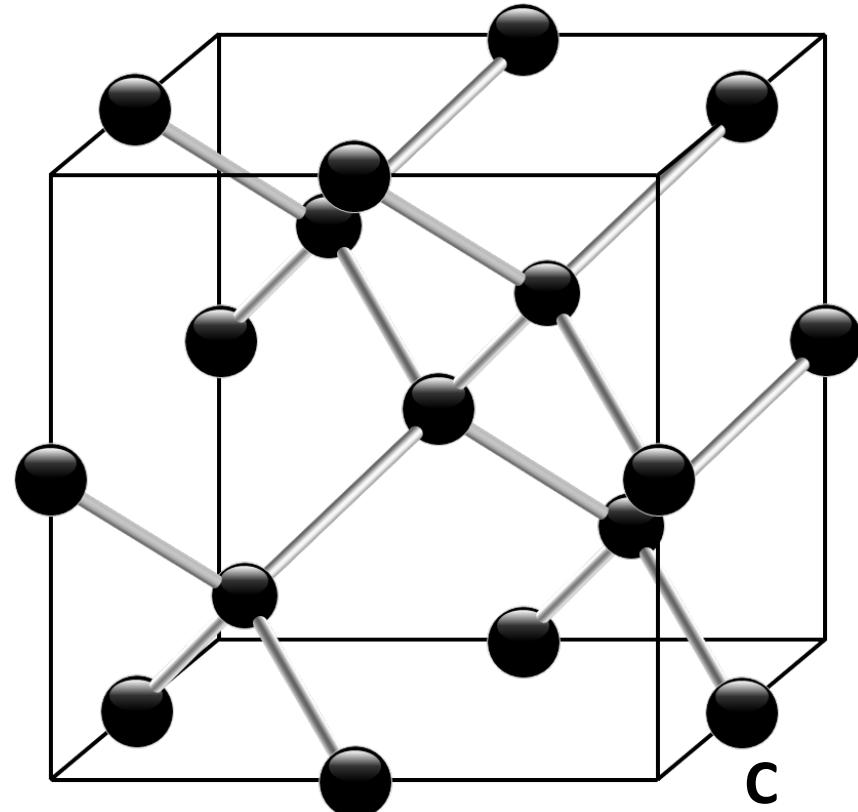
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# Diamond envy



©GIA



$$\rho_N = 1.77 \times 10^{23} \text{ cm}^{-3}$$

# Diamond NV

**Synthetic (CVD) diamond**

$2^2 \times 0.5 \text{ mm}^3$ , \$700 (E6)

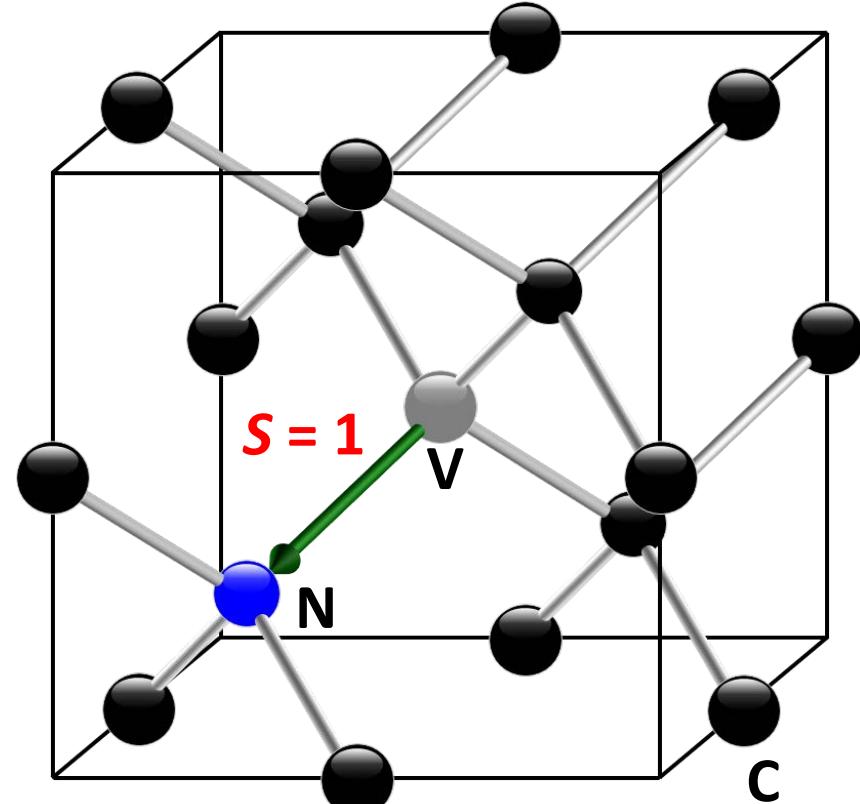
[N] < 5 ppb, [NV] < 0.03 ppb



*Not like...*



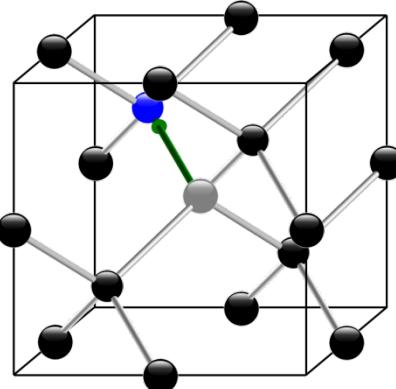
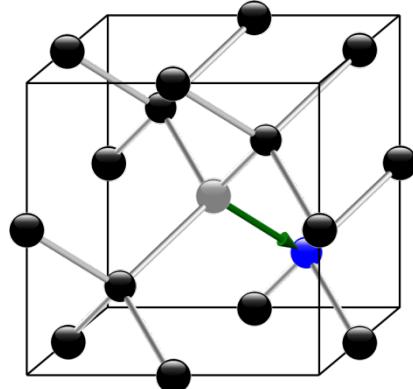
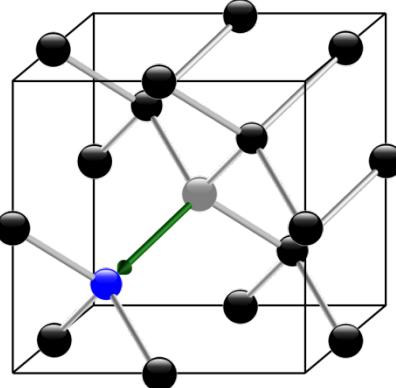
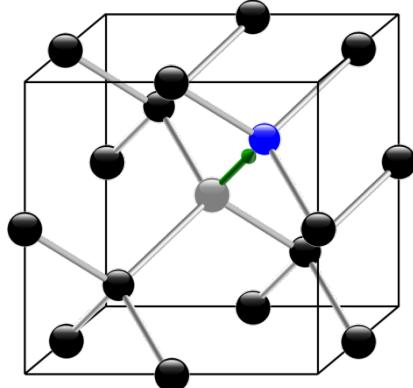
©GIA



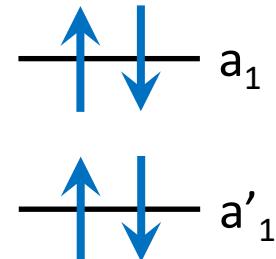
$$\rho_N = 1.77 \times 10^{23} \text{ cm}^{-3}$$

# Crystal & energy level structures

- Negatively-charged ( $\text{NV}^-$ )
- 4  $sp^3$  orbitals, 6  $e^-$  (5 from the defect, 1 captured)
- $C_{3v}$  (symmetry axis = quantization axis)



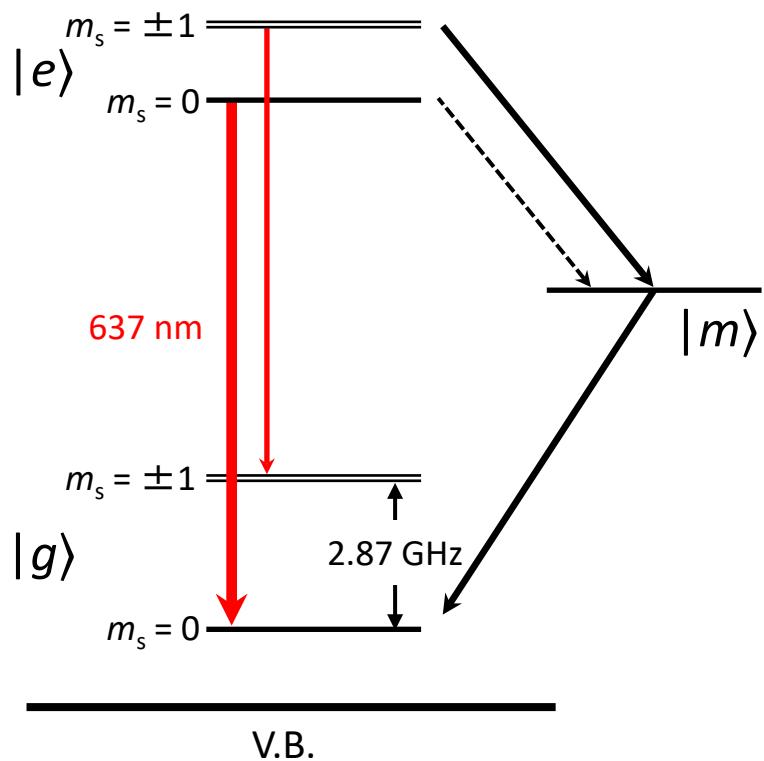
Effective spin-1 system  
( $e^2$ -hole spin-triplet)



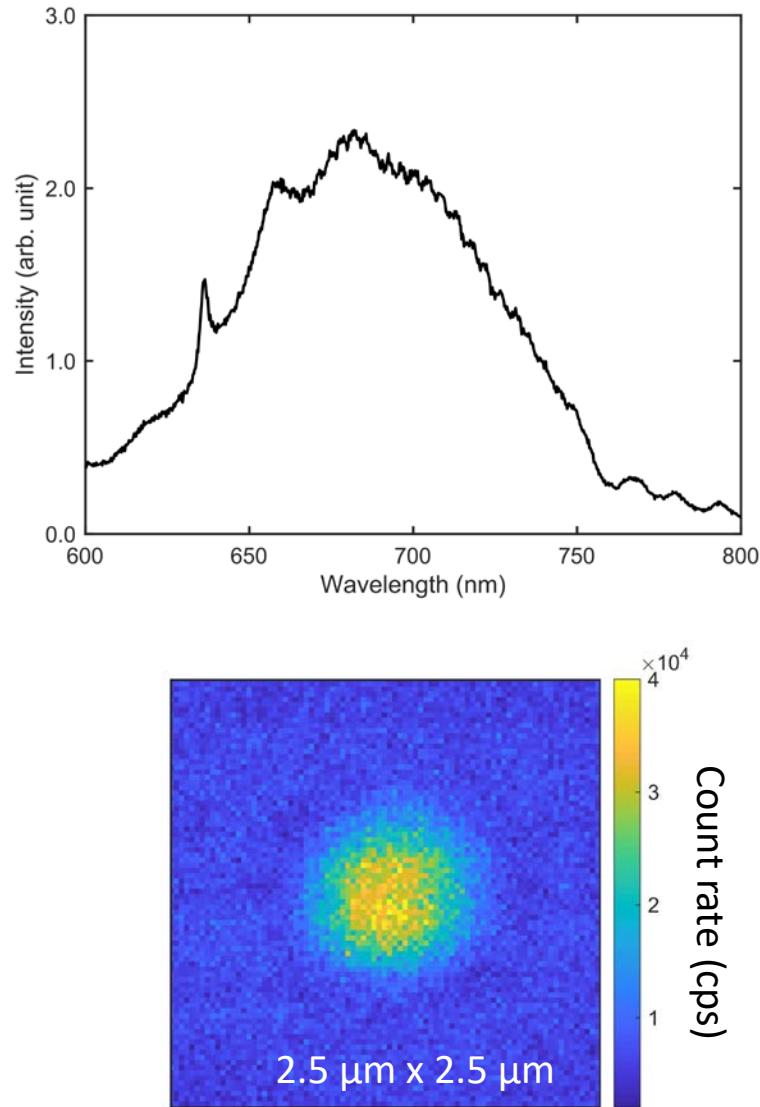
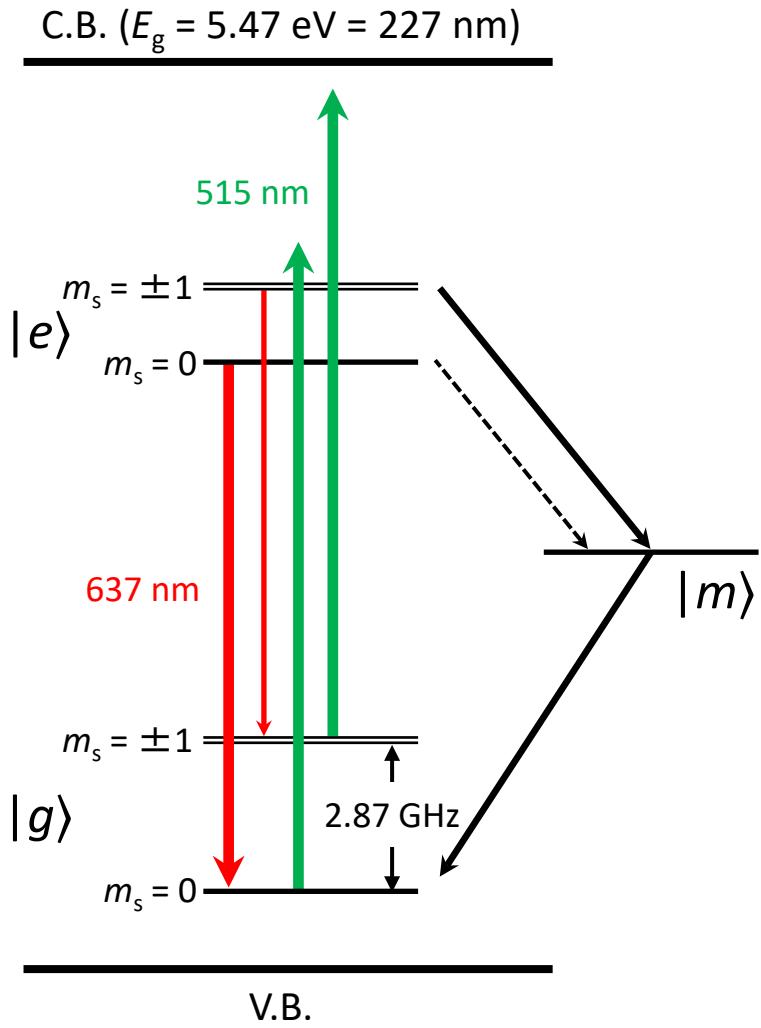
# Energy levels

C.B. ( $E_g = 5.47 \text{ eV} = 227 \text{ nm}$ )

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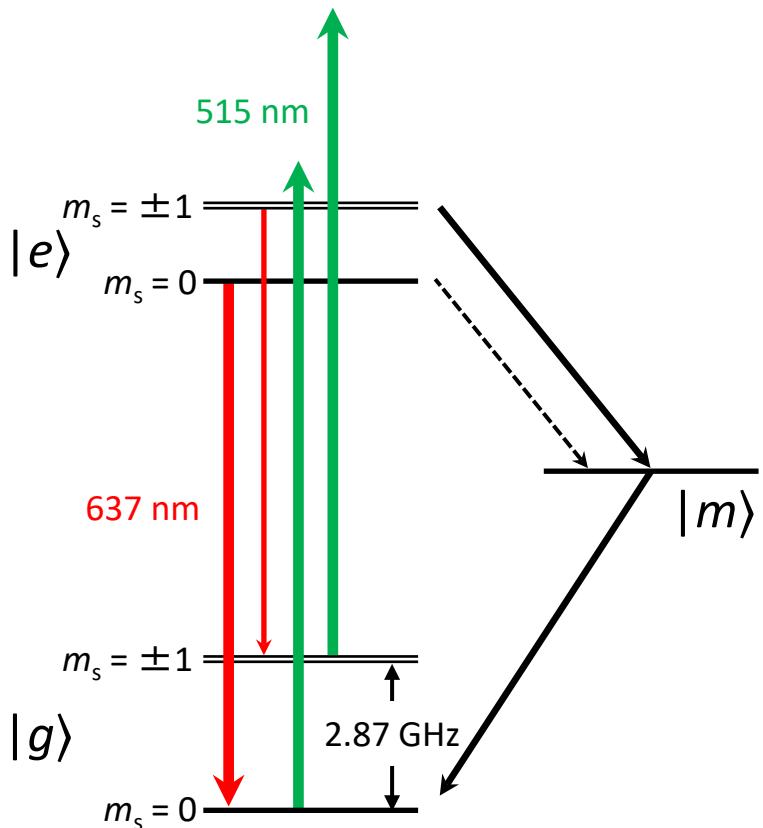


# PL spectroscopy & imaging

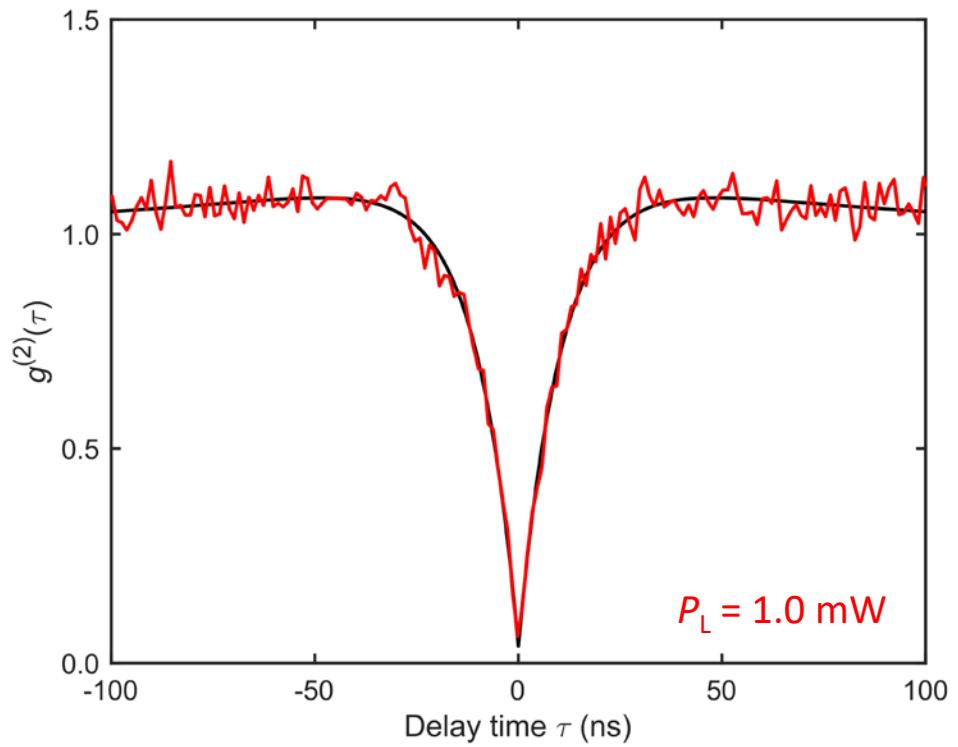


# Photon statistics

C.B. ( $E_g = 5.47 \text{ eV} = 227 \text{ nm}$ )



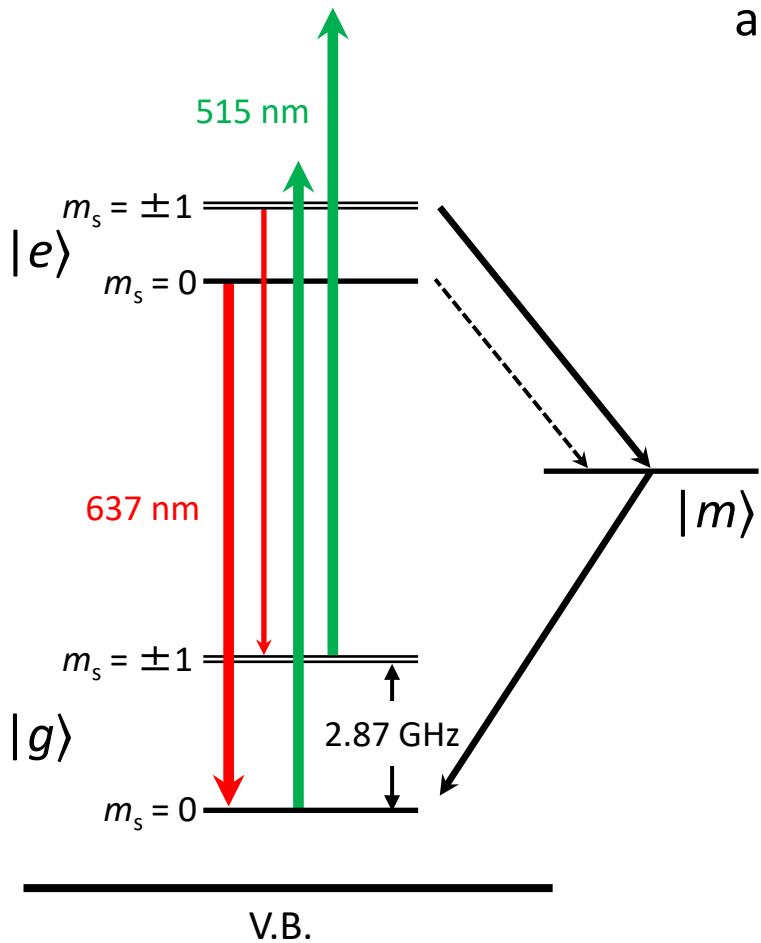
One photon at a time



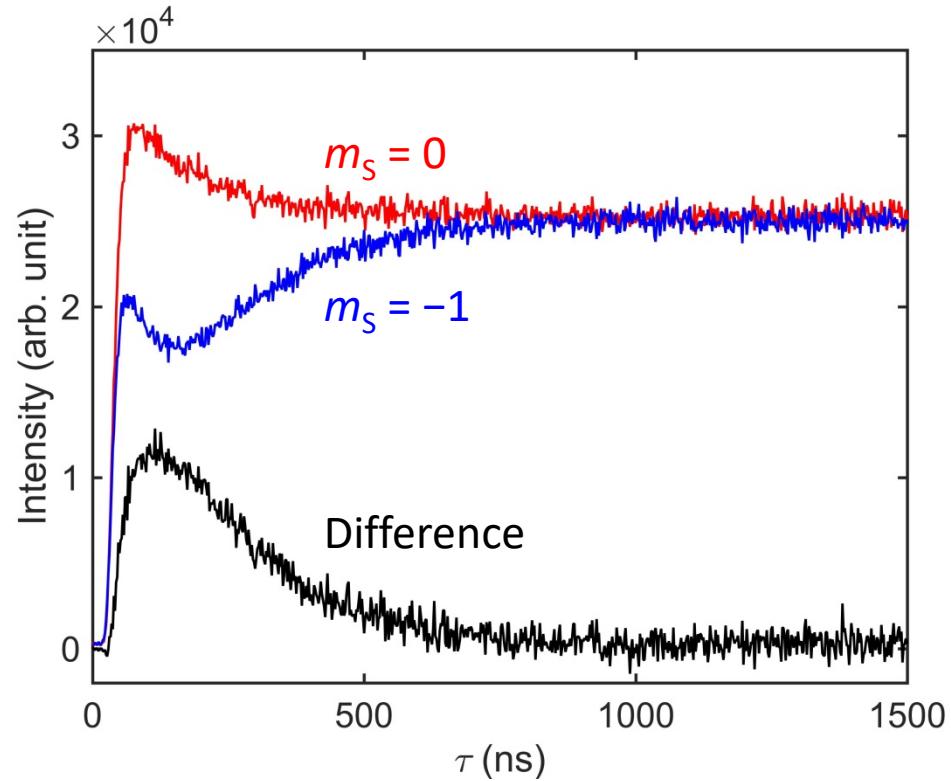
V.B.

# Time-resolved fluorescence

C.B. ( $E_g = 5.47 \text{ eV} = 227 \text{ nm}$ )

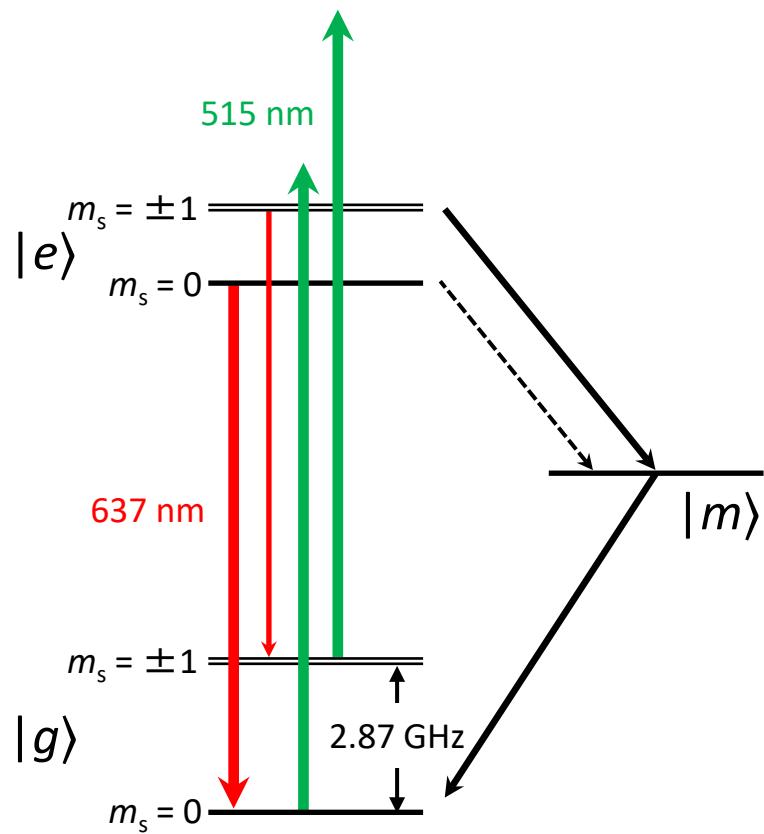


The **non-radiative & spin-selective** channel provides a means to **read out & initialize** the NV spin



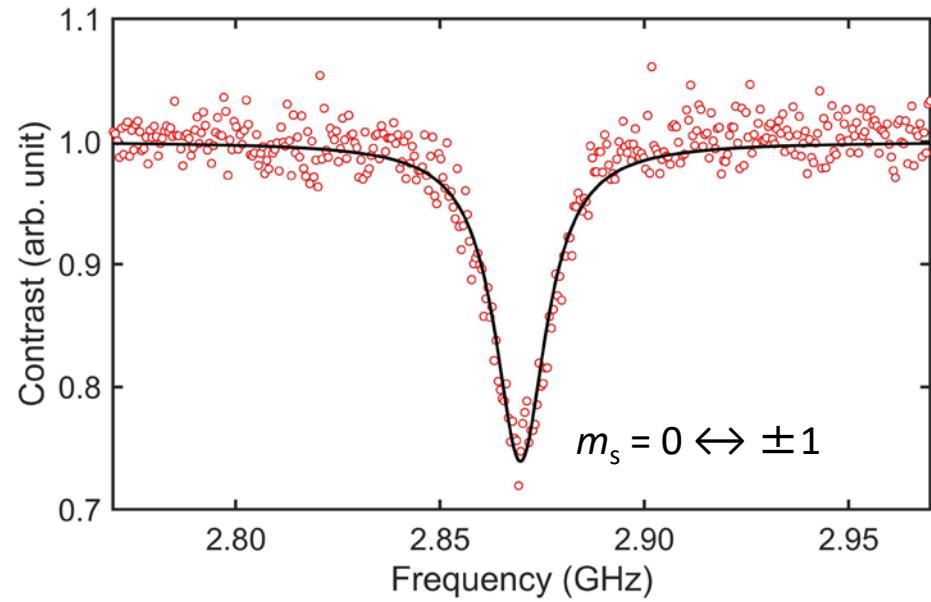
# CW ODMR at $B_0 = 0$

C.B. ( $E_g = 5.47$  eV = 227 nm)



**Zero-field splitting**  $H = DS_z^2$

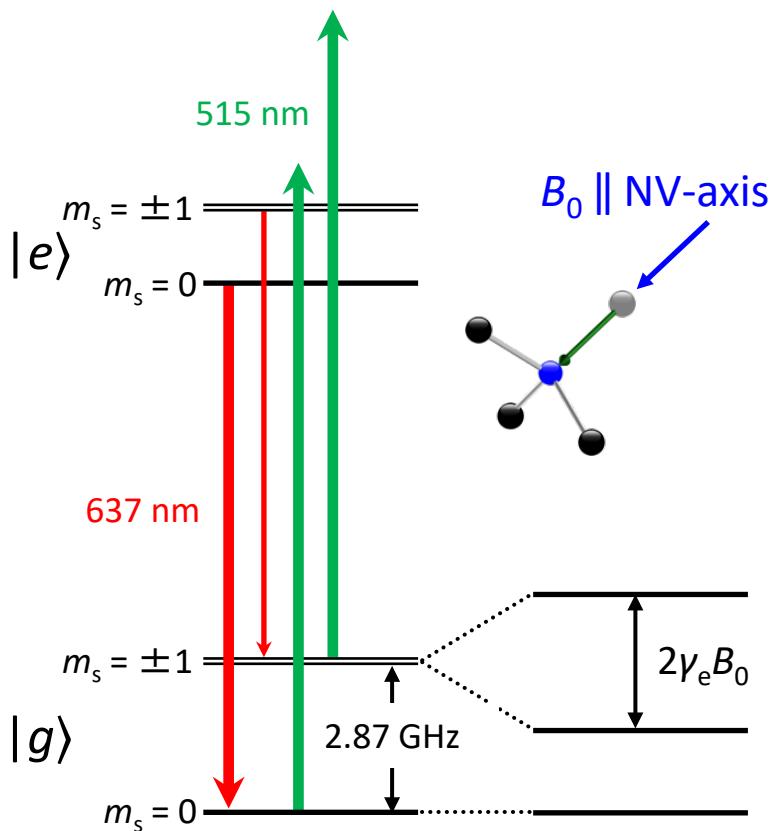
$$D = 2.87 \text{ GHz}$$



V.B.

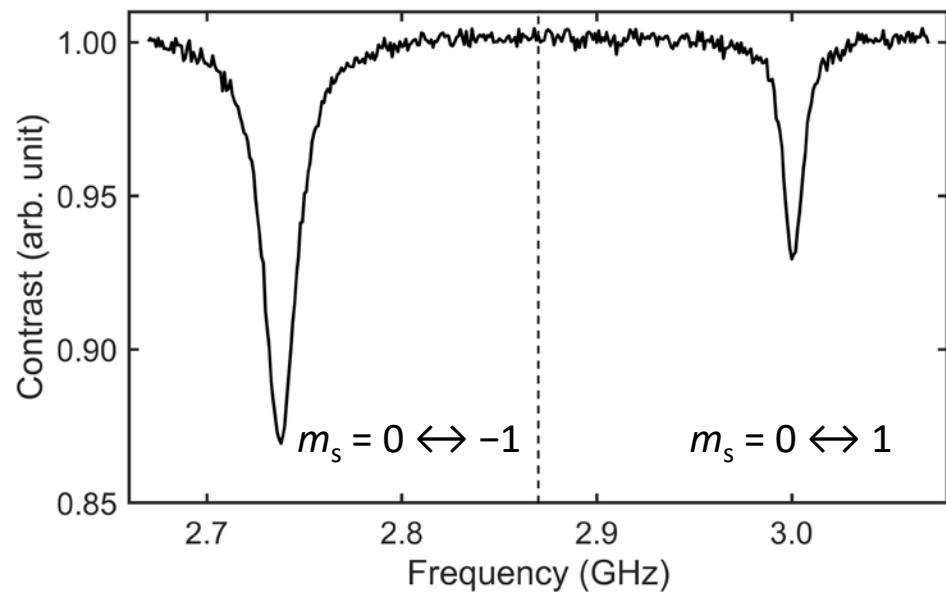
# CW ODMR at $B_0 > 0$

C.B. ( $E_g = 5.47 \text{ eV} = 227 \text{ nm}$ )



$$\text{Zeeman} \quad H = DS_z^2 + \gamma_e B_0 S_z$$

$$\gamma_e = 28 \text{ MHz/mT}$$



$$B_0 = 4.7 \text{ mT} (2.87 \pm 0.132 \text{ GHz})$$

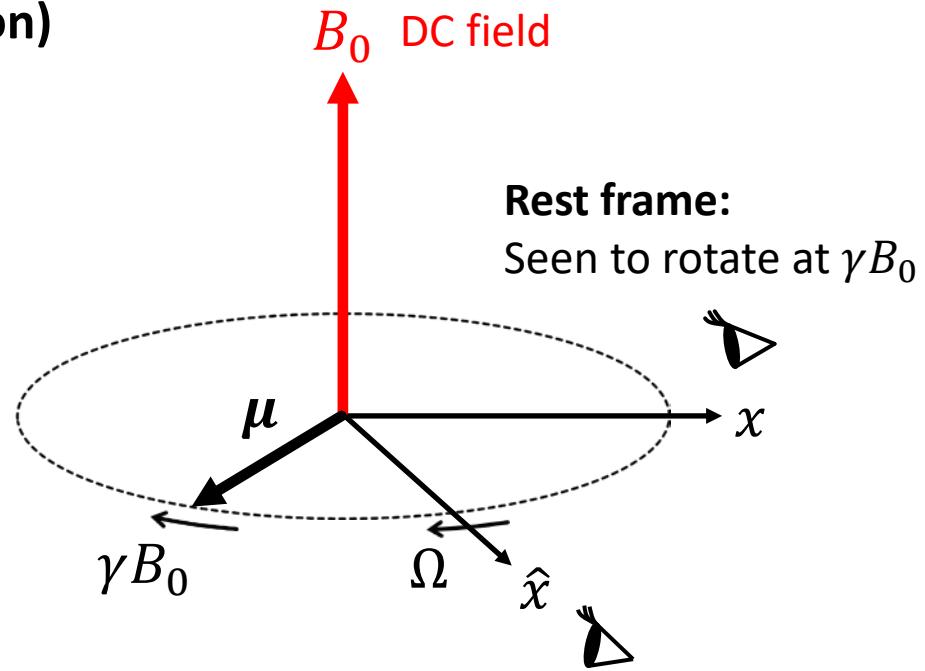
# Magnetic resonance

Torque equation (Larmor precession)

$$\frac{d\mu}{dt} = \mu \times \gamma B_0$$

↑  
↑  
Gyromagnetic ratio

Magnetic moment:  $\mu = \gamma J$

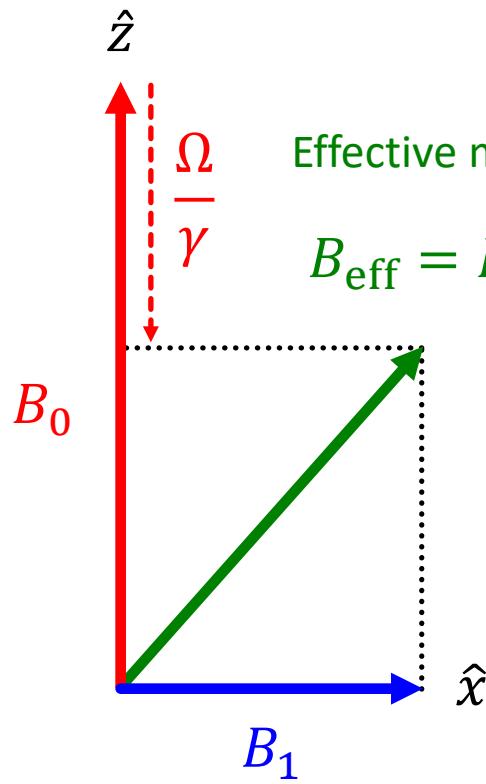


Frame rotating at angular velocity  $\Omega$ :  
Rotate slower...why?



DC field along the  $z$  direction becomes weaker

# Magnetic resonance

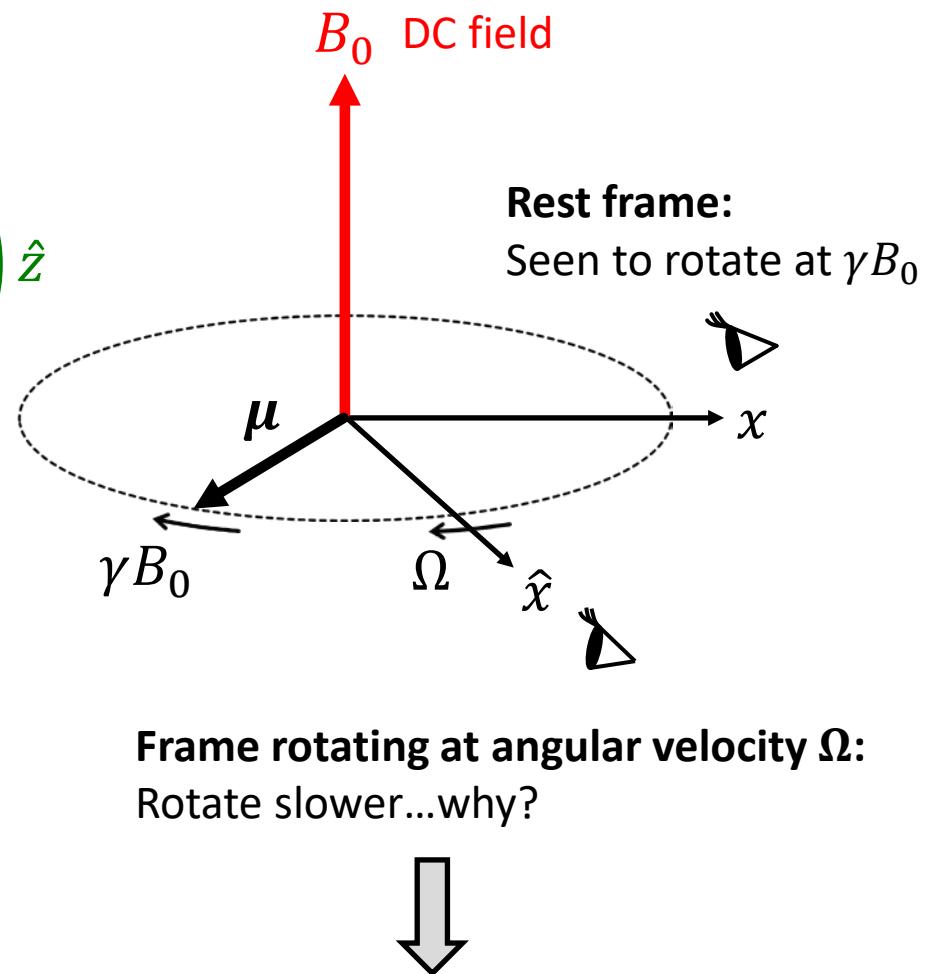


Effective magnetic field

$$B_{\text{eff}} = B_1 \hat{x} + \left( B_0 - \frac{\Omega}{\gamma} \right) \hat{z}$$

$B_1$

AC field rotating in the  $xy$  plane at  $\Omega$



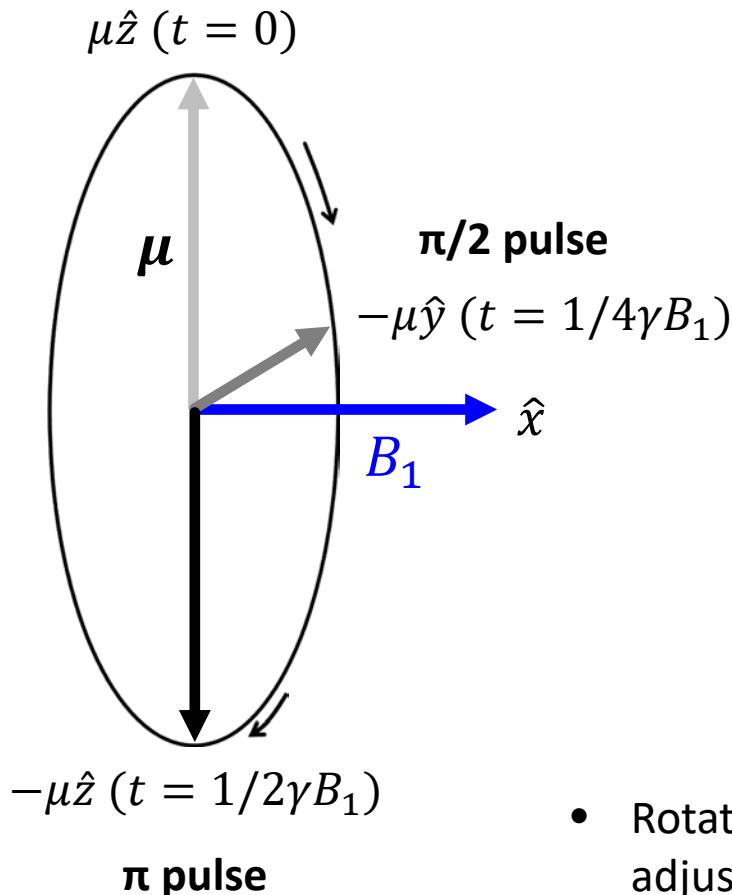
Frame rotating at angular velocity  $\Omega$ :  
Rotate slower...why?



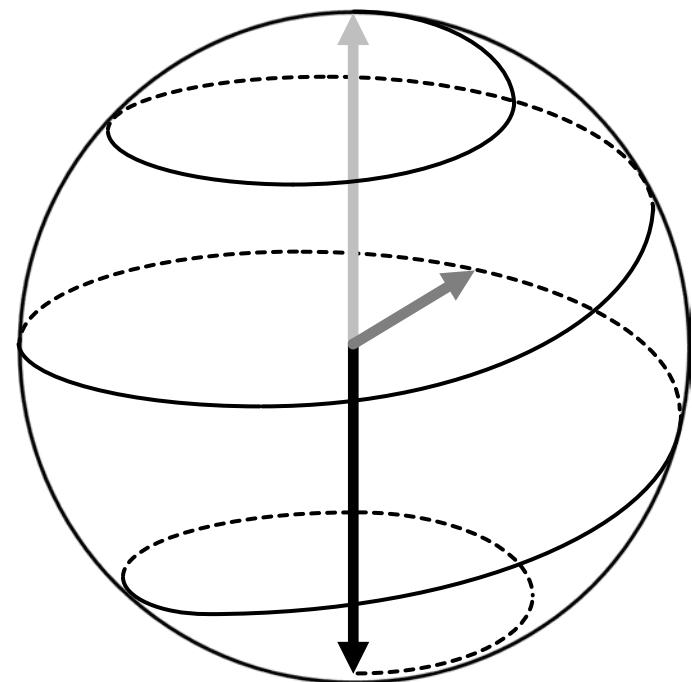
DC field along the  $z$  direction becomes weaker

# Magnetic resonance

Frame rotating at  $\Omega = \gamma B_0$



Rest (non-resonant) frame



- Rotations about the  $\pm \hat{x}, \pm \hat{y}$  axes are realized by adjusting the microwave phases
- Rotation about the  $\hat{z}$  axis is superposed when observed from the rest (non-resonant) frame

# Quantum bit

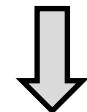
**Qubit, spin-1/2 (NV is spin-1!)**

$$\left\{ \begin{array}{l} |"0"\rangle \equiv |m_s = 0\rangle \\ |"1"\rangle \equiv |m_s = -1\rangle \end{array} \right.$$

**Superposition state**

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

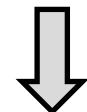
$$|\alpha|^2 + |\beta|^2 = 1$$



$$|\psi\rangle = e^{i\gamma} \left( \cos \frac{\theta}{2} |0\rangle + e^{i\phi} \sin \frac{\theta}{2} |1\rangle \right)$$

$$\begin{aligned} 0 &\leq \theta \leq \pi \\ 0 &\leq \gamma, \phi < 2\pi \end{aligned}$$

Global phase



$$|\psi\rangle = \cos \frac{\theta}{2} |0\rangle + e^{i\phi} \sin \frac{\theta}{2} |1\rangle$$

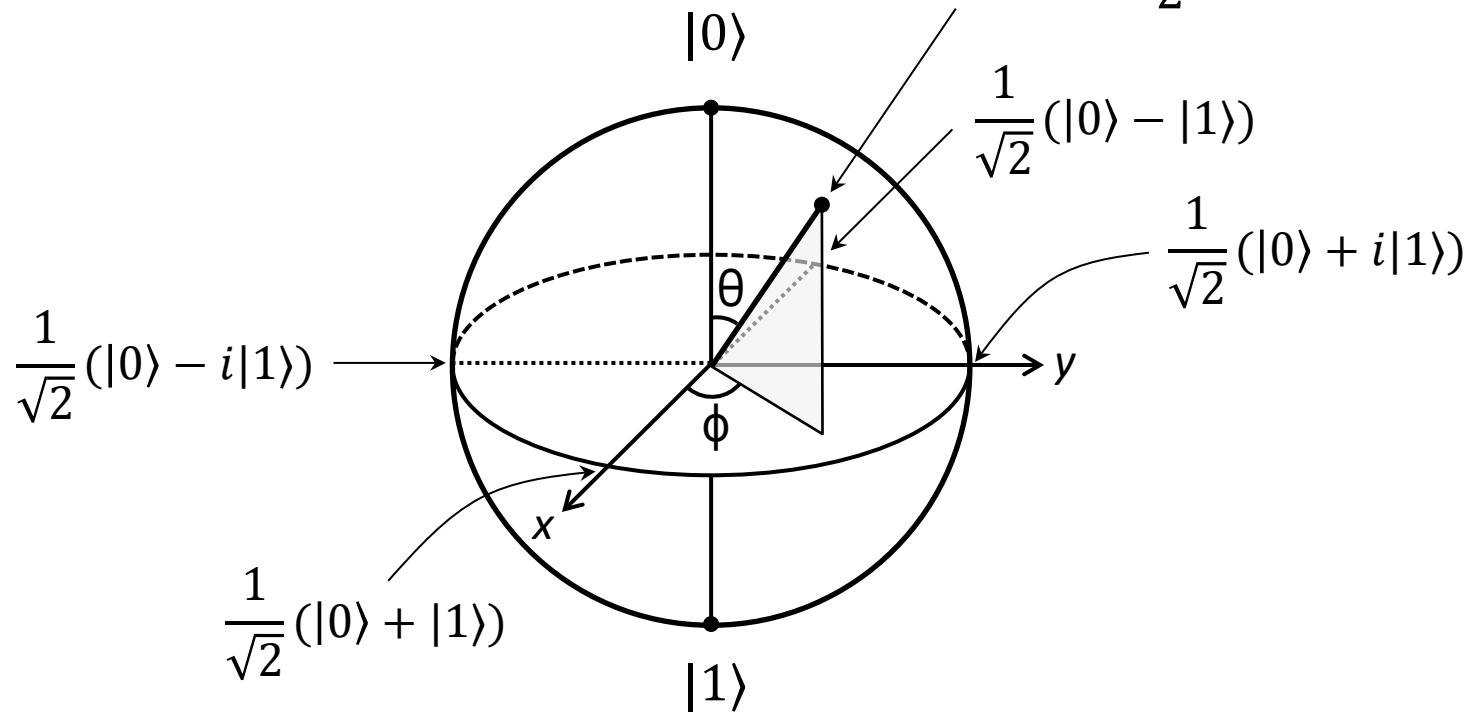
$$\begin{aligned} 0 &\leq \theta \leq \pi \\ 0 &\leq \phi < 2\pi \end{aligned}$$

# Bloch sphere

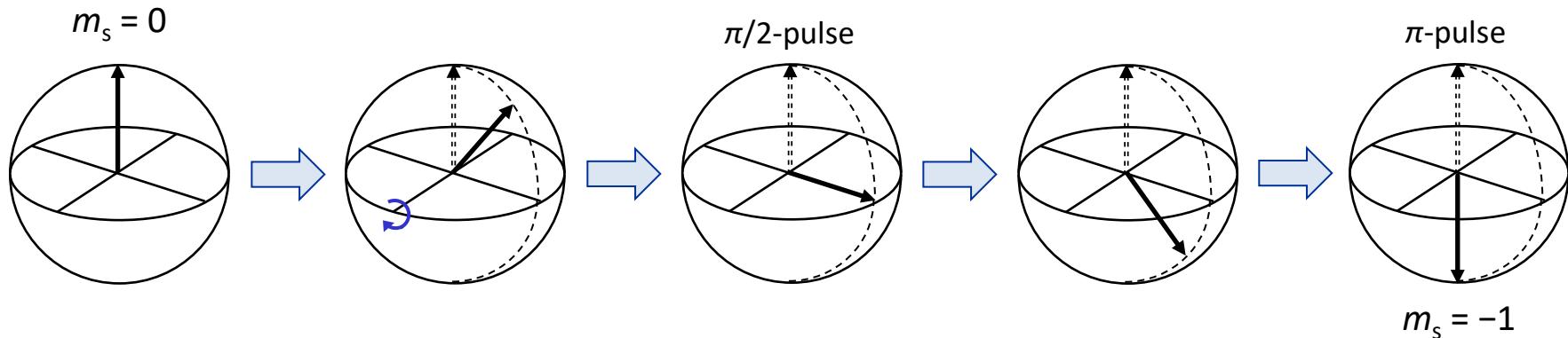
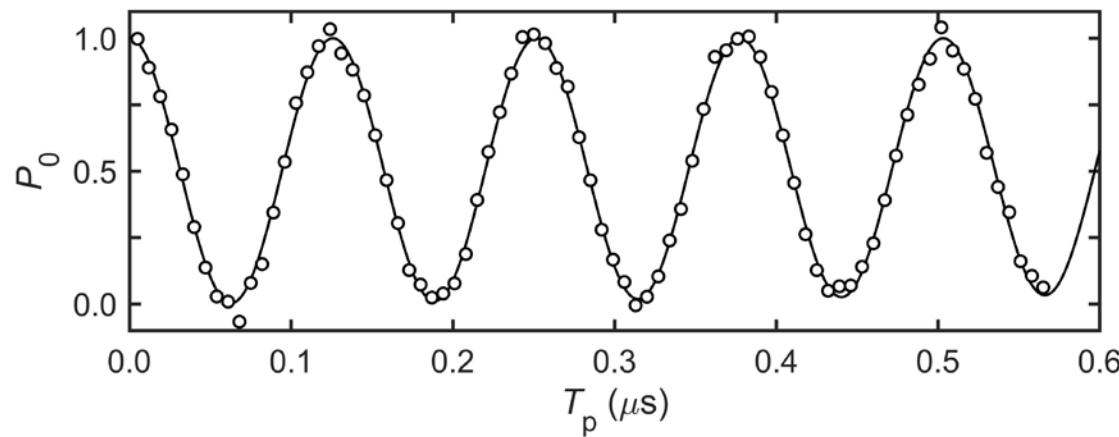
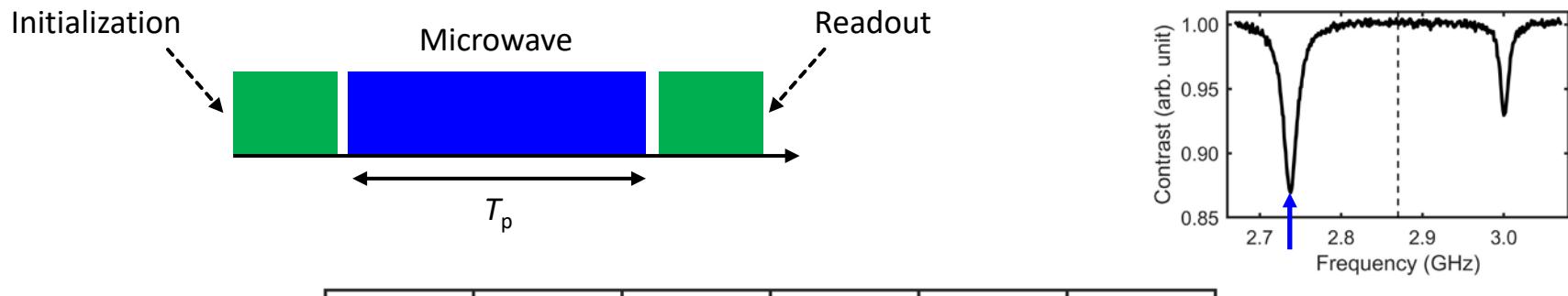
Qubit, spin-1/2 (NV is spin-1!)

$$\begin{cases} \text{"0"} \rangle \equiv |m_s = 0\rangle \\ \text{"1"} \rangle \equiv |m_s = -1\rangle \end{cases}$$

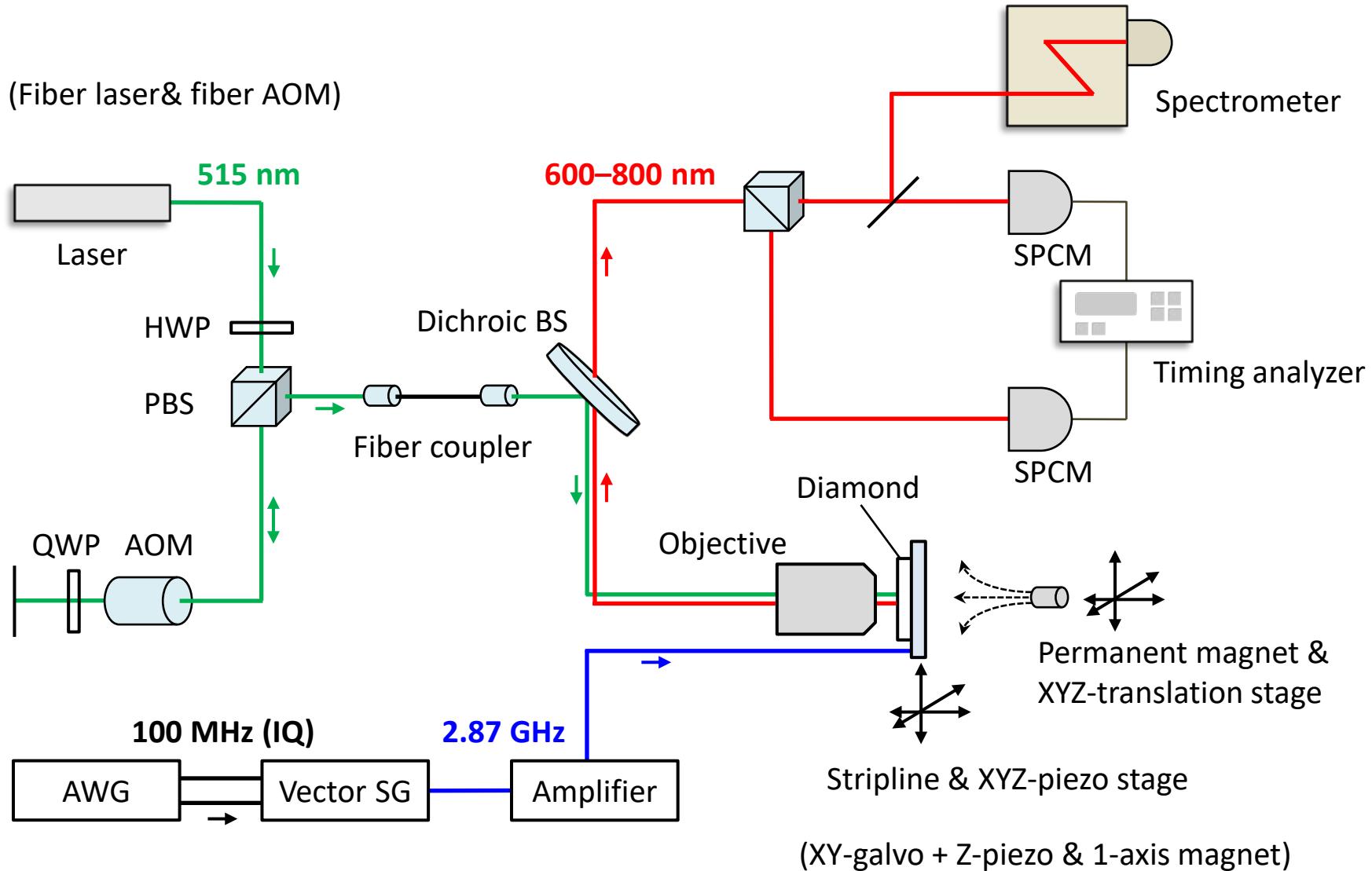
$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\phi}\sin\frac{\theta}{2}|1\rangle$$



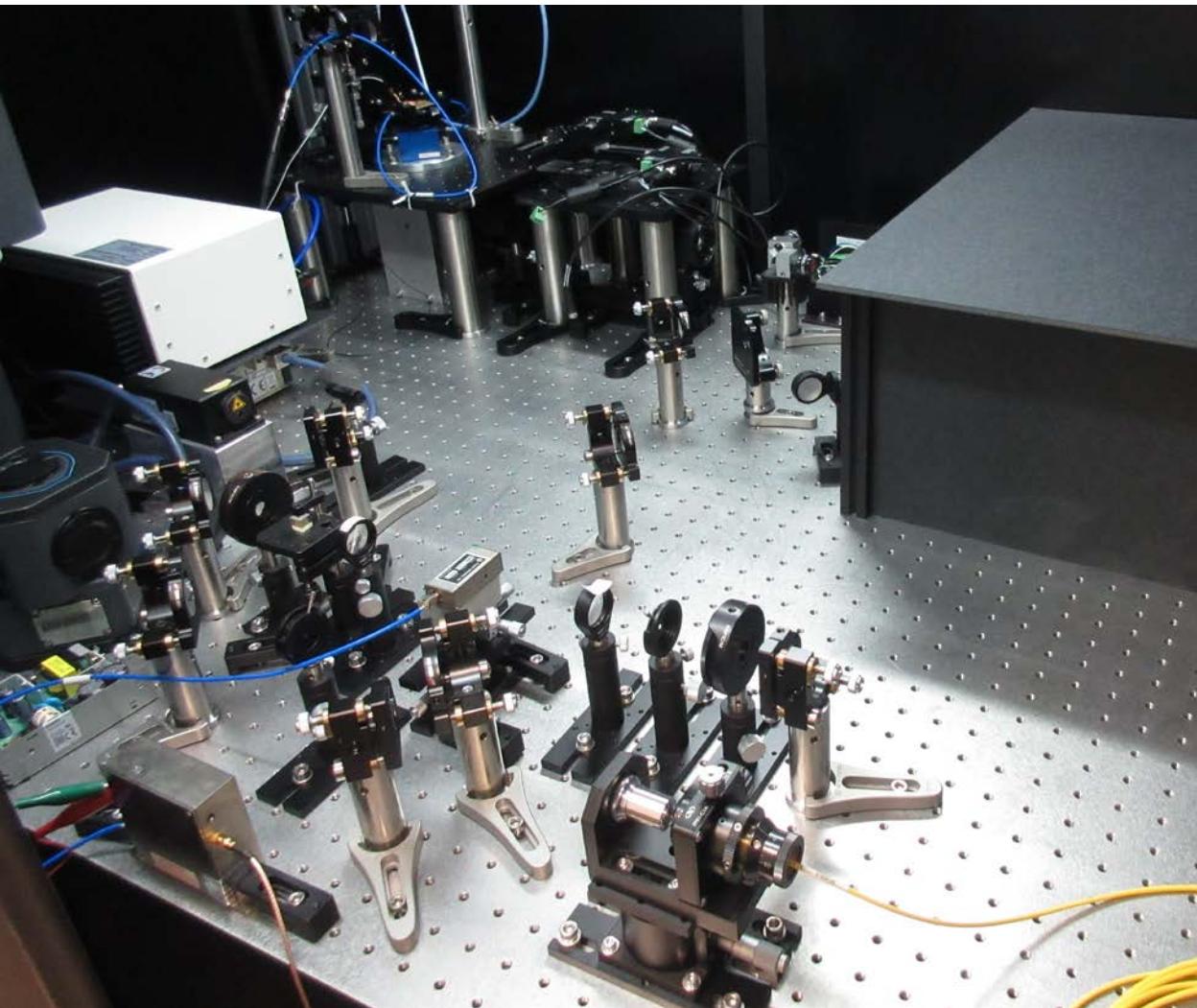
# Rabi oscillation



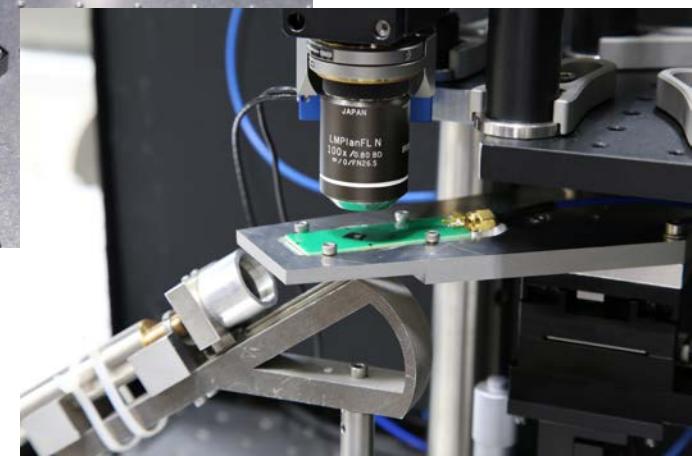
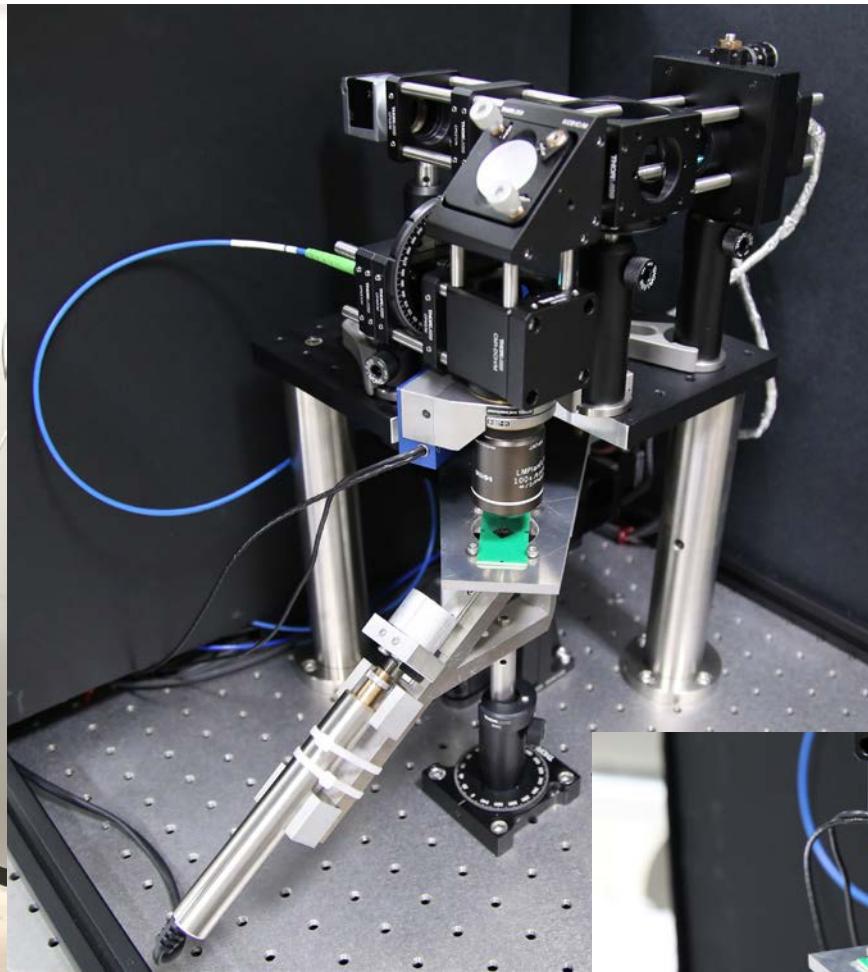
# Experimental setup



# Experimental setup



# Experimental setup

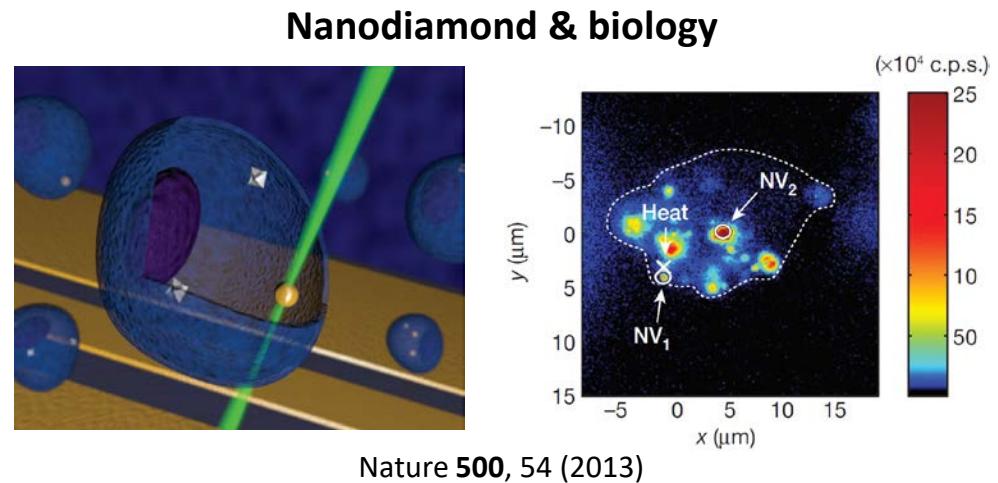


# Outline

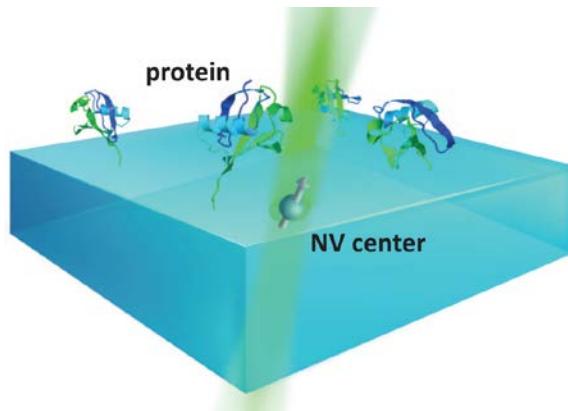
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  - Structure
  - Optical properties
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- **Quantum sensing**
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  - Detection and localization of a single  $^{13}\text{C}$  nuclear spin
  - Ultrahigh resolution sensing

# Quantum sensing with NV centers

- $B, E, T, S\dots$
- DC & AC modes
- Wide temperature range
- Nondestructive
- High spatial resolution
- Various modalities

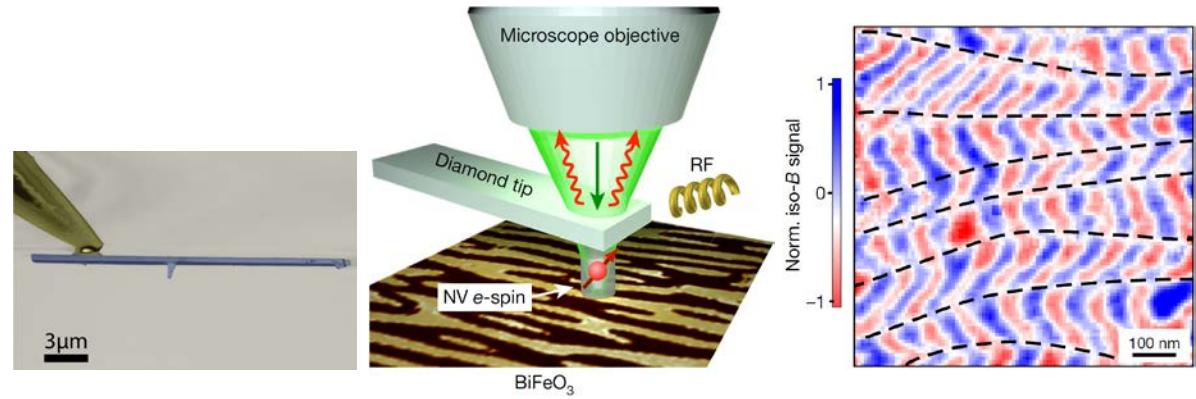


## Near-surface NV center & NMR



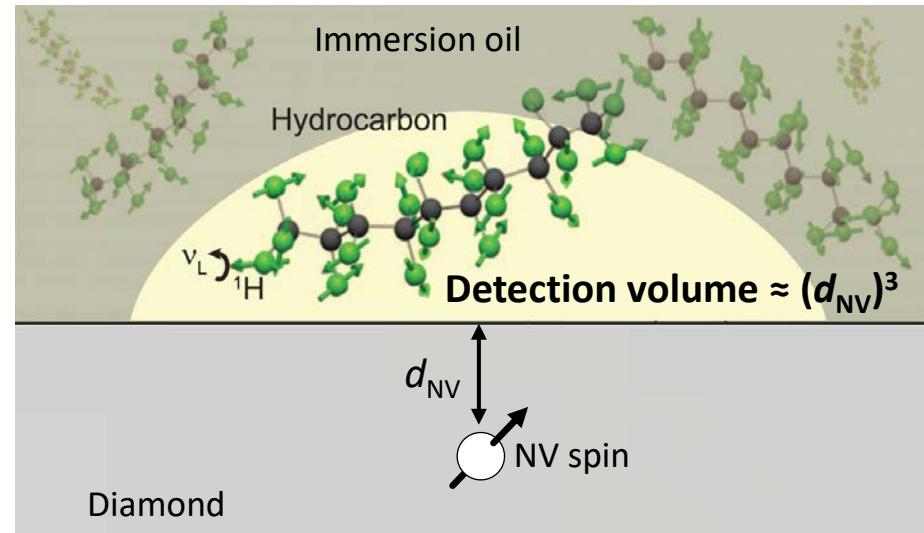
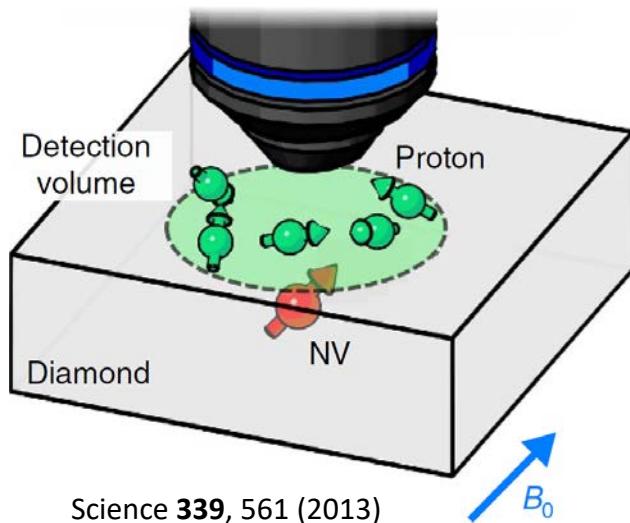
Science 351, 836 (2016)

## Scanning probe & condensed matter



Rev. Sci. Instrum. 87, 063703 (2016); Nature 549, 252 (2017)

# Nuclear spin sensing

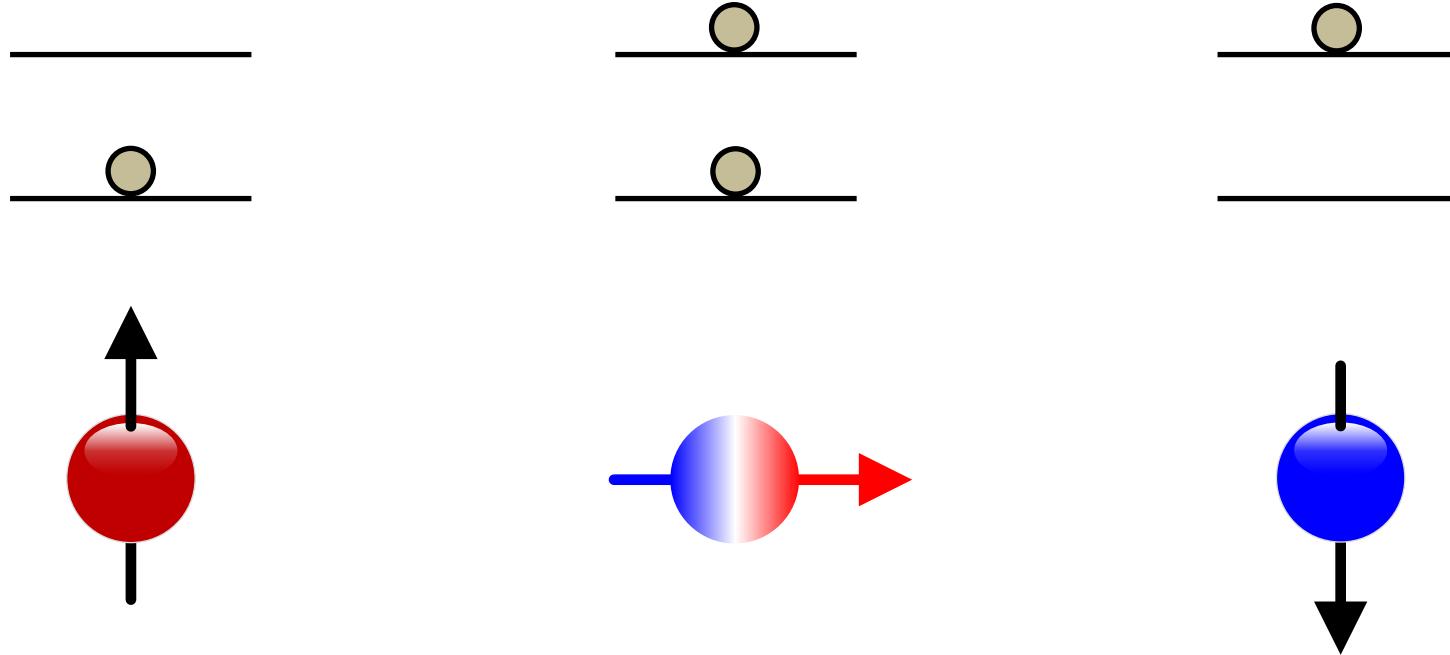


Nuclear spins **precess** at  $f_{ac} = \text{a few kHz--MHz}$  under  $B_0$

→ **Weak AC magnetic field** on the NV spin

→ Detect using **quantum coherence**

# Quantum coherence



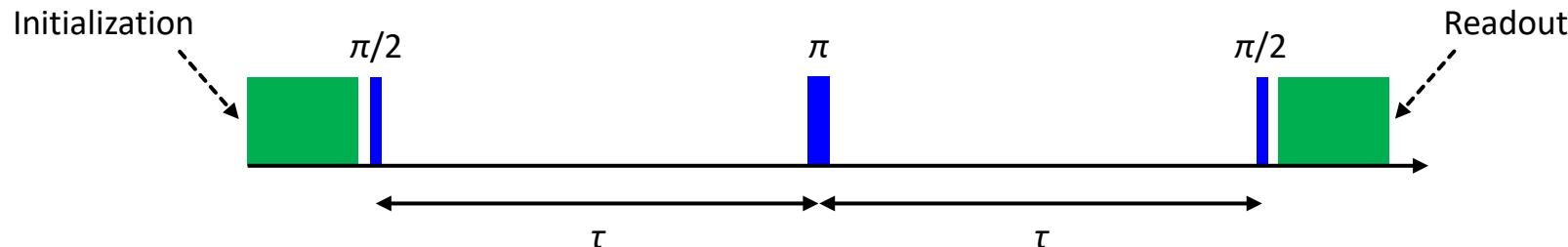
$$|0\rangle \equiv |m_s = 0\rangle$$

$$|\Psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

$$|1\rangle \equiv |m_s = -1\rangle$$

**$T_2$ : measure of how long a superposition state is preserved**

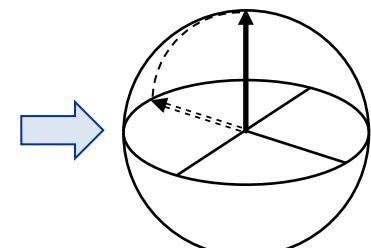
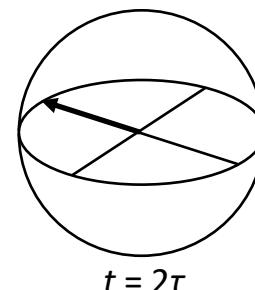
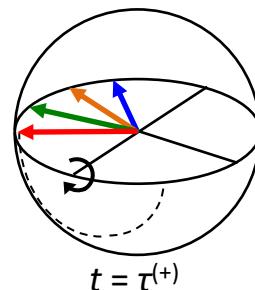
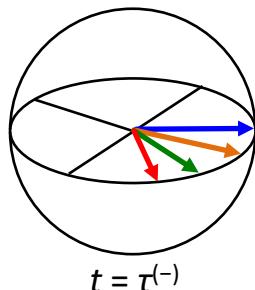
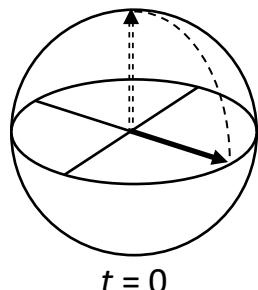
# Spin echo



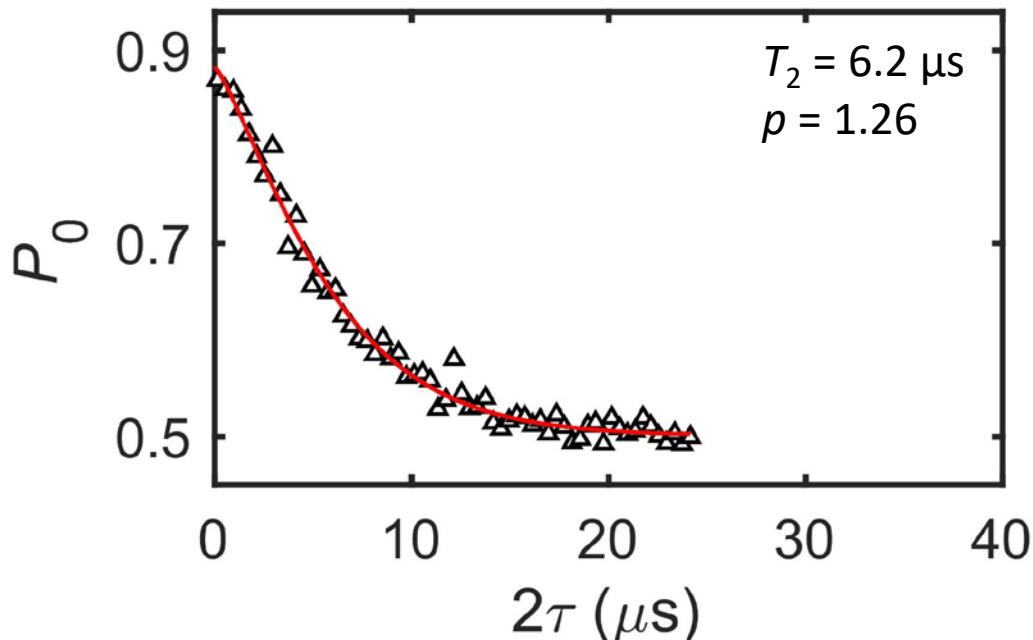
**Modulation function**



**Phase accumulation  
by DC field**



# Coherence time



**Stretched exponential decay**

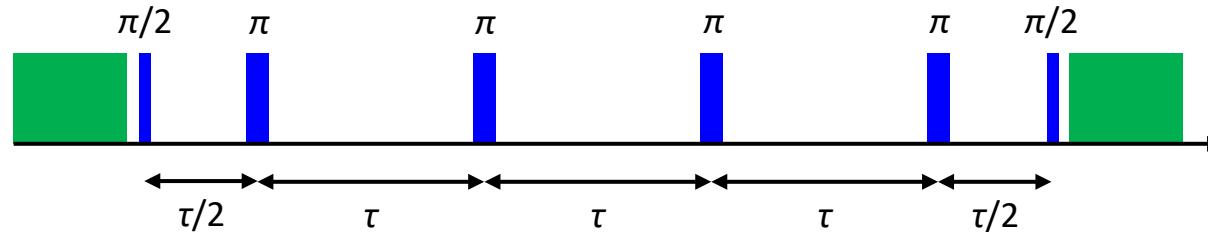
$$\exp \left[ - \left( \frac{2\tau}{T_2} \right)^p \right]$$

**Near-surface NV center**

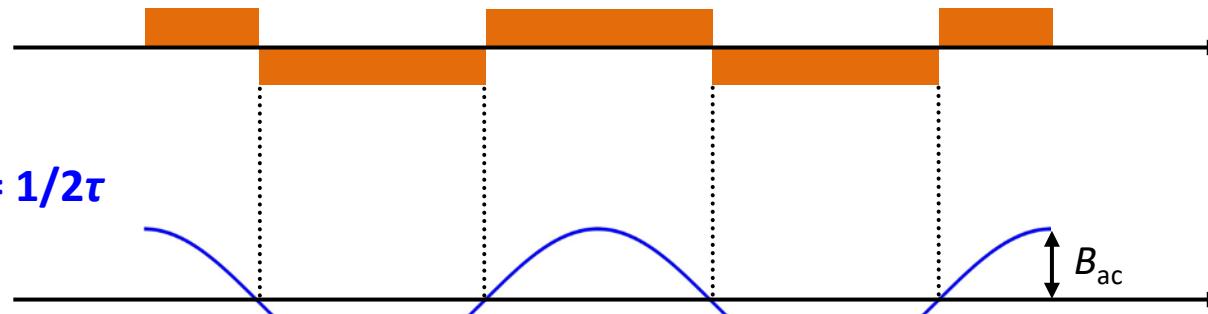
- N<sup>+</sup> implantation into <sup>12</sup>C ( $I = 0$ ) layer
- $d_{\text{NV}} = 6.26 \text{ nm}$
- $B_0 = 23.5 \text{ mT}$

# AC magnetometry

CP ( $N = 4$ )

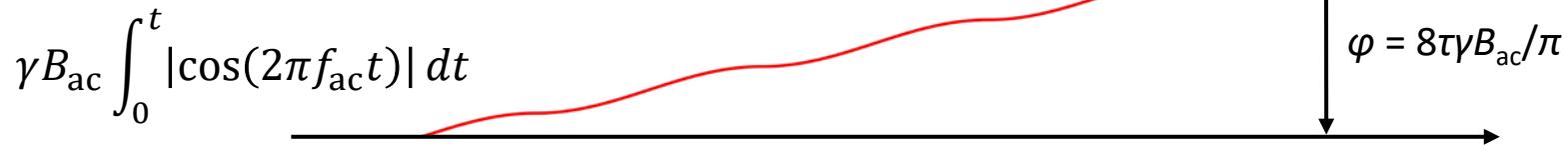


Modulation function

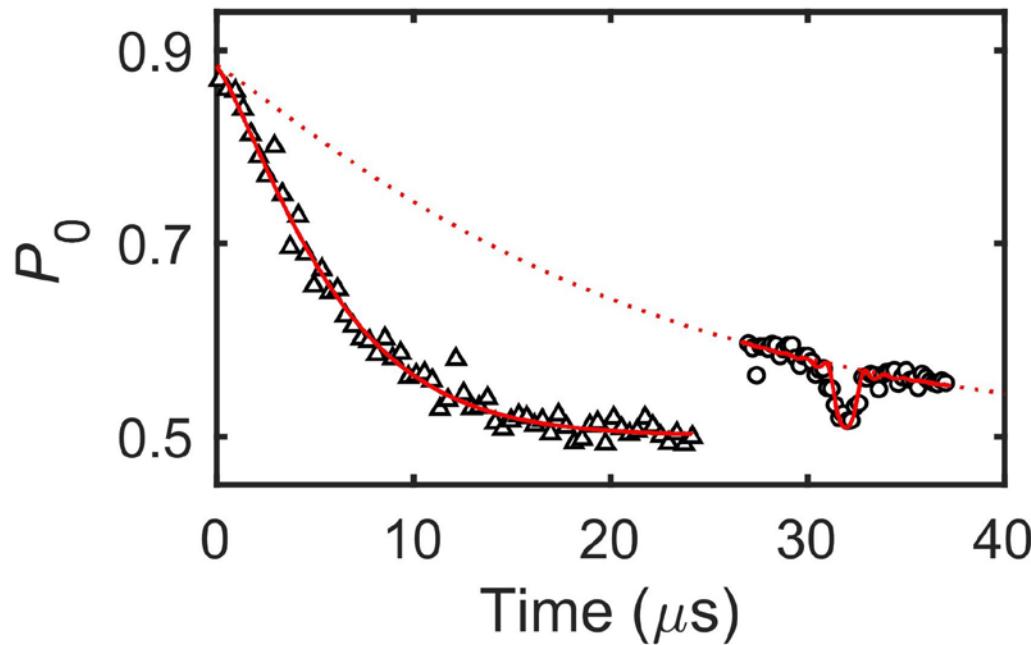


AC field at  $f_{ac} = 1/2\tau$

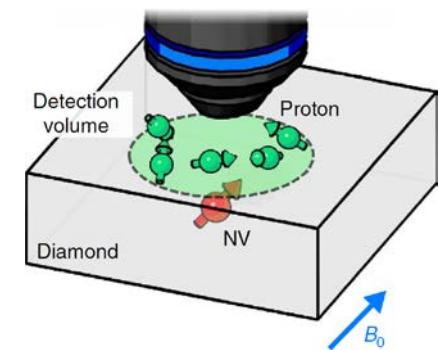
*Sensor phase buildup = loss of coherence*



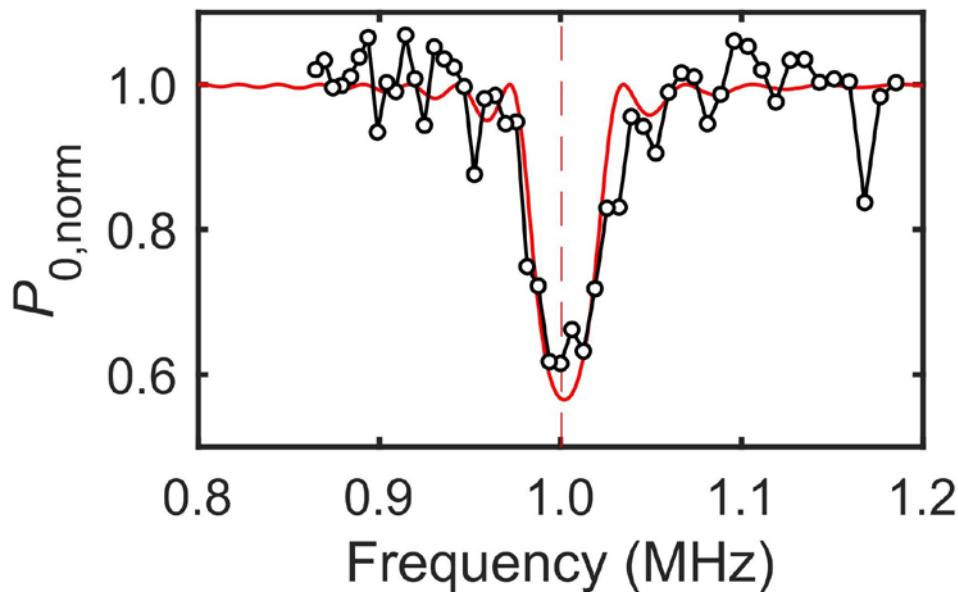
# Nuclear spin sensing



- $T_2 = 6.2 \mu\text{s} @ B_0 = 23.5 \text{ mT}$
- $N = 64$  (XY16)
- $(2\tau)^{-1} = 64/(2 \times 32 \mu\text{s}) = 1 \mu\text{s}$   
     $\rightarrow \gamma_H B_0 = (42.577 \text{ kHz/mT}) \times B_0 = 1.00 \text{ MHz}$



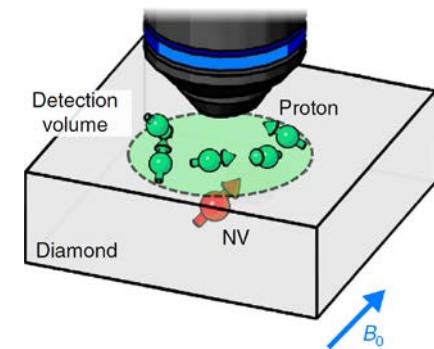
# Nuclear spin sensing



$$C(\tau) = f(B_{\text{rms}})$$

$$B_{\text{rms}} = \frac{\mu_0}{4\pi} h \gamma_H \sqrt{\frac{5\pi\rho}{96d_{\text{NV}}^3}}$$

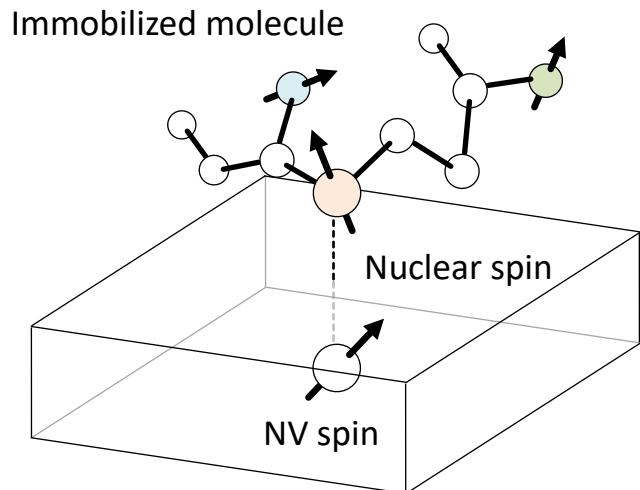
Phys. Rev. B **93**, 045425 (2016)



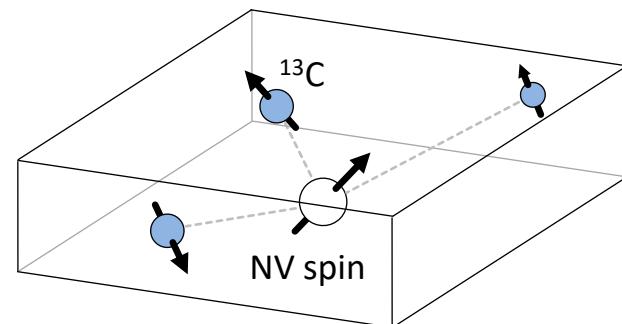
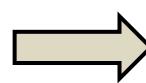
- Proton density  $\rho = 6 \times 10^{28} \text{ m}^{-3}$  (known)
- $d_{\text{NV}} = 6.26 \text{ nm}$
- $B_{\text{rms}} \approx 560 \text{ nT}$
- Detection volume  $(d_{\text{NV}})^3 \approx 0.25 \text{ zL}$  (zepto =  $10^{-21}$ )
- # of proton  $\rho(d_{\text{NV}})^3 \approx 1500$
- Thermal pol. ( $10^{-7}$ ) vs. statistical pol.  $(1500)^{0.5} \approx 39$

# Toward single-molecular imaging

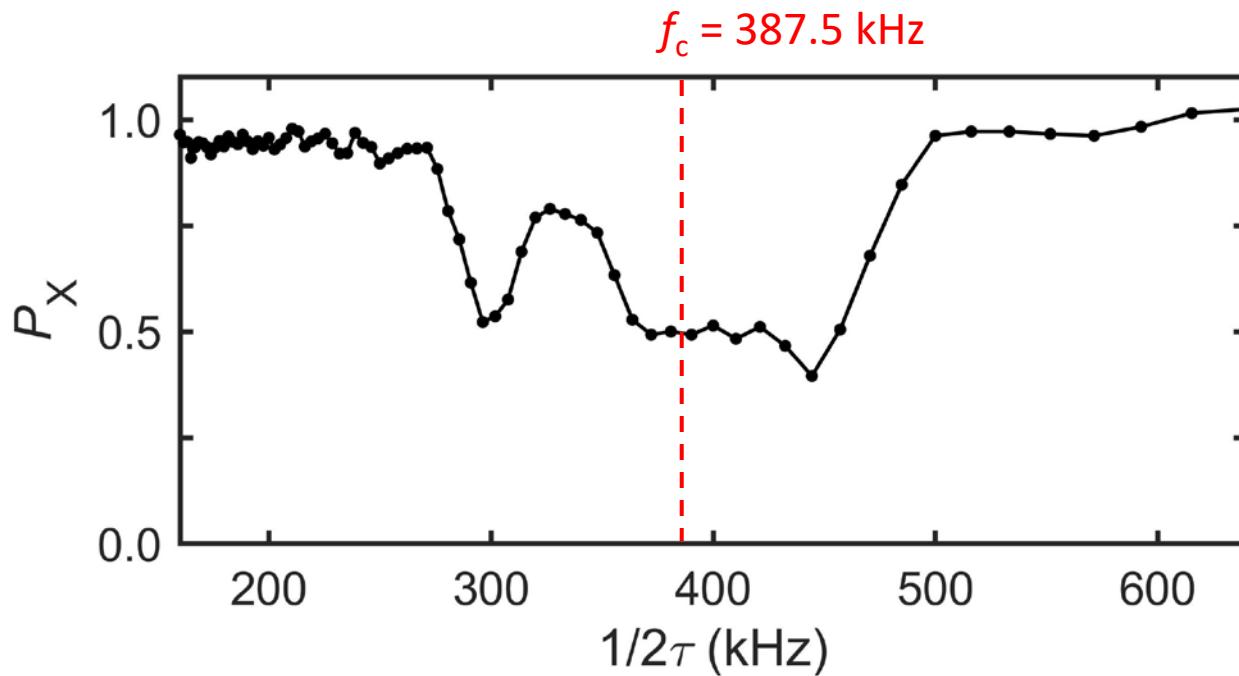
- **Strategy**
  - Detect **individual nuclear spins** contained in a single molecule
  - Determine their **nuclear species (& chemical shifts)** and **positions**
- **Practical issues**
  - Preparation of high-quality near-surface NV centers
  - Accurate positioning of single molecules/proteins near the sensor



Use  $^{13}\text{C}$  (1.1%) in diamond as a testbed



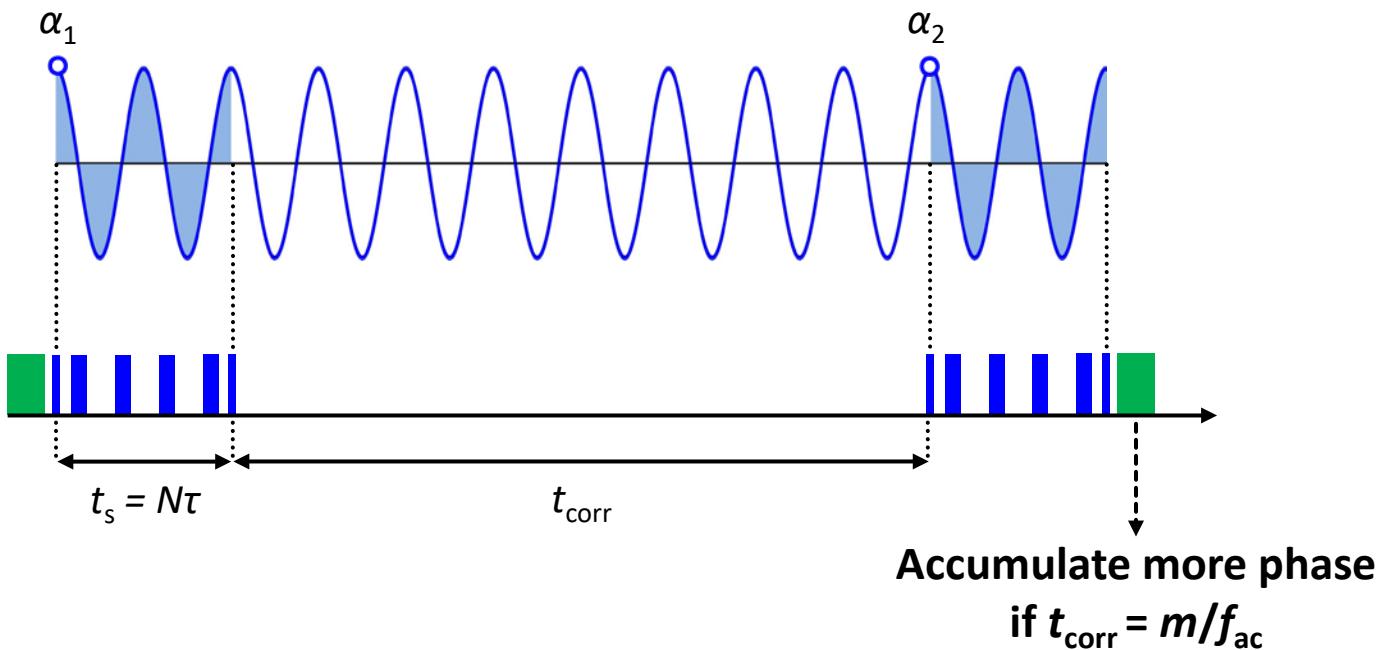
# Nuclear spin sensing



- Single NV in bulk ( $[^{13}\text{C}] = 1.1\%$ ,  $d_{\text{NV}} \approx 50 \mu\text{m}$ )
- $N = 16$
- $f_c = \gamma_c B_0 = 10.705 \text{ kHz/mT} \times 36.2 \text{ mT}$

# Correlation spectroscopy

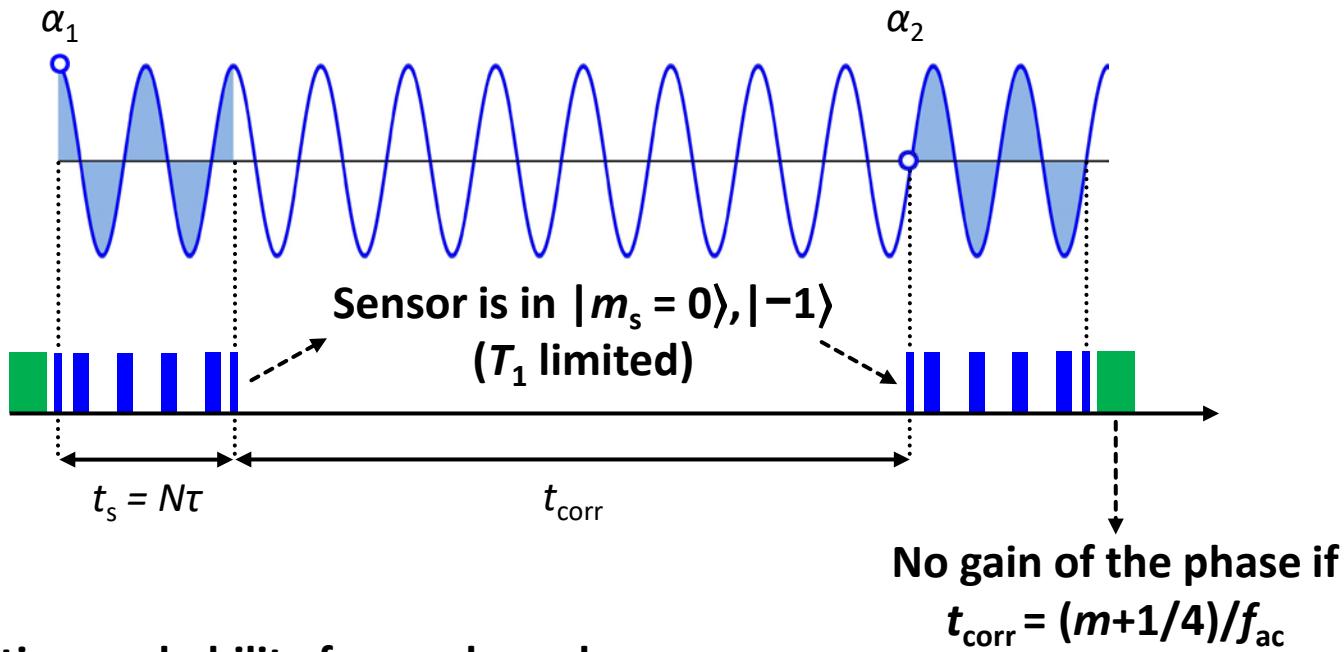
AC field at  $f_{ac}$



- Nature Commun. **4**, 1651 (2013) Laraoui *et al.*  
Phys. Rev. Appl. **4**, 024004 (2015) Kong *et al.*  
Nature Commun. **6**, 8527 (2015) Staudacher *et al.*  
Phys. Rev. Lett. **116**, 197601 (2016) Boss *et al.*

# Correlation spectroscopy

AC field at  $f_{ac}$



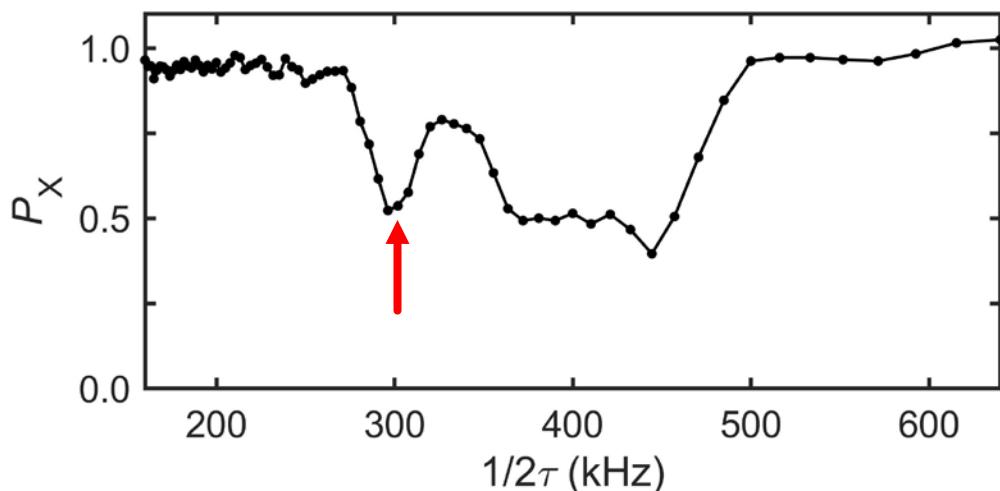
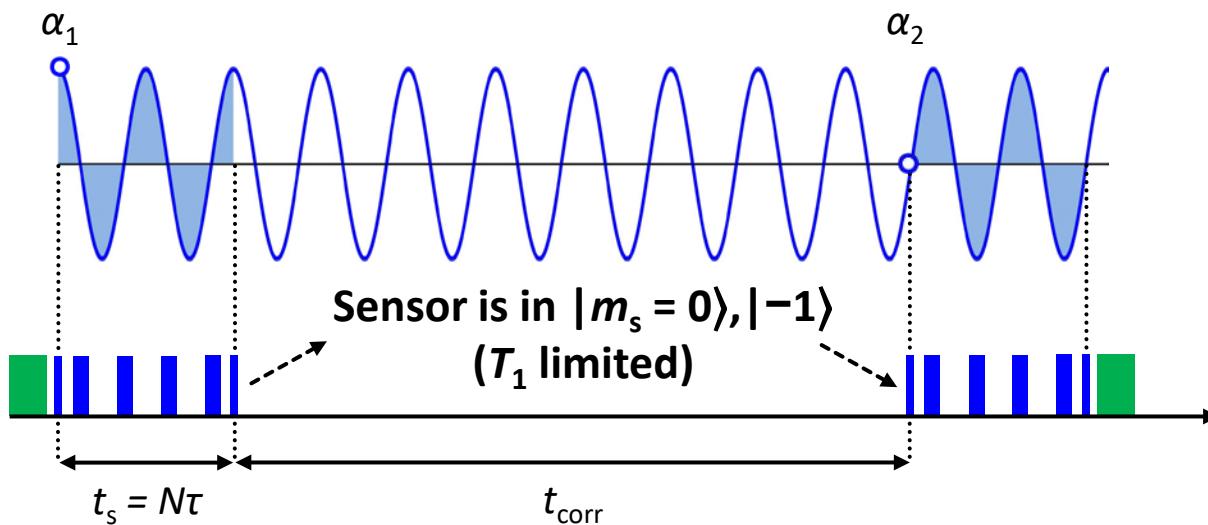
The transition probability for random phases

$$p(t_1) \approx \frac{1}{2} \left\{ 1 - \frac{1}{2} \left( \frac{\gamma B_{ac} t_s}{\pi} \right)^2 \cos(2\pi f_{ac} t_{corr}) \right\}$$

- Nature Commun. **4**, 1651 (2013) Laraoui *et al.*  
Phys. Rev. Appl. **4**, 024004 (2015) Kong *et al.*  
Nature Commun. **6**, 8527 (2015) Staudacher *et al.*  
Phys. Rev. Lett. **116**, 197601 (2016) Boss *et al.*

# Correlation spectroscopy

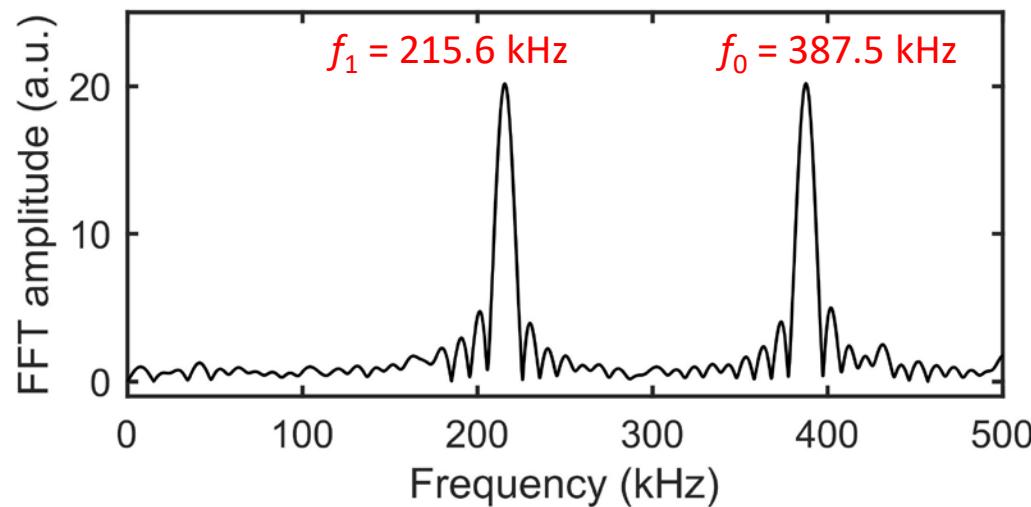
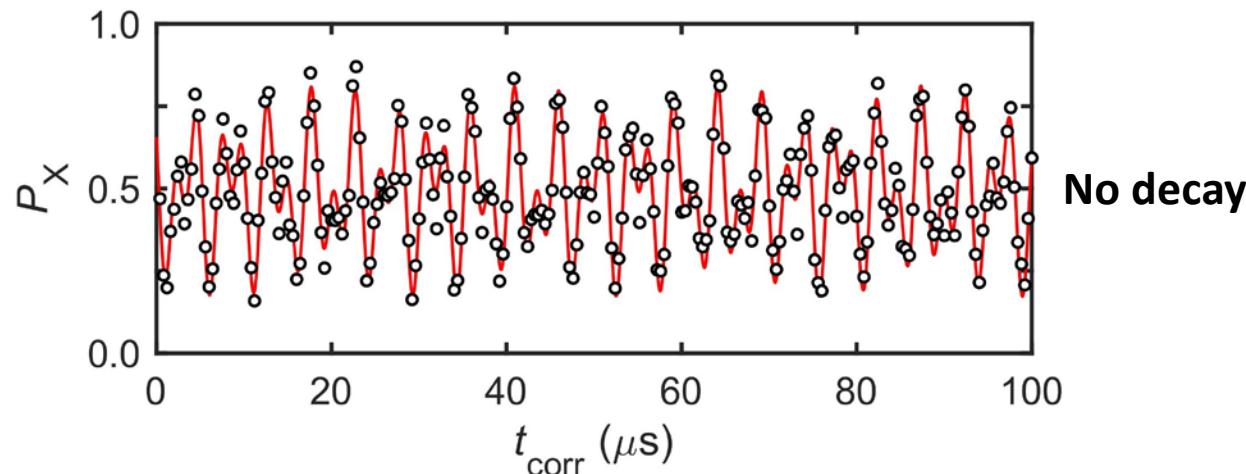
AC field at  $f_{ac}$



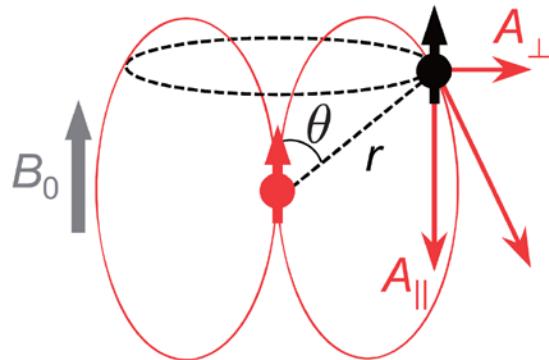
Where to look at?

- $f_t = 1/2\tau = 301.6$  kHz
- $\tau = 1.7875$   $\mu$ s

# Correlation spectroscopy of a nucleus



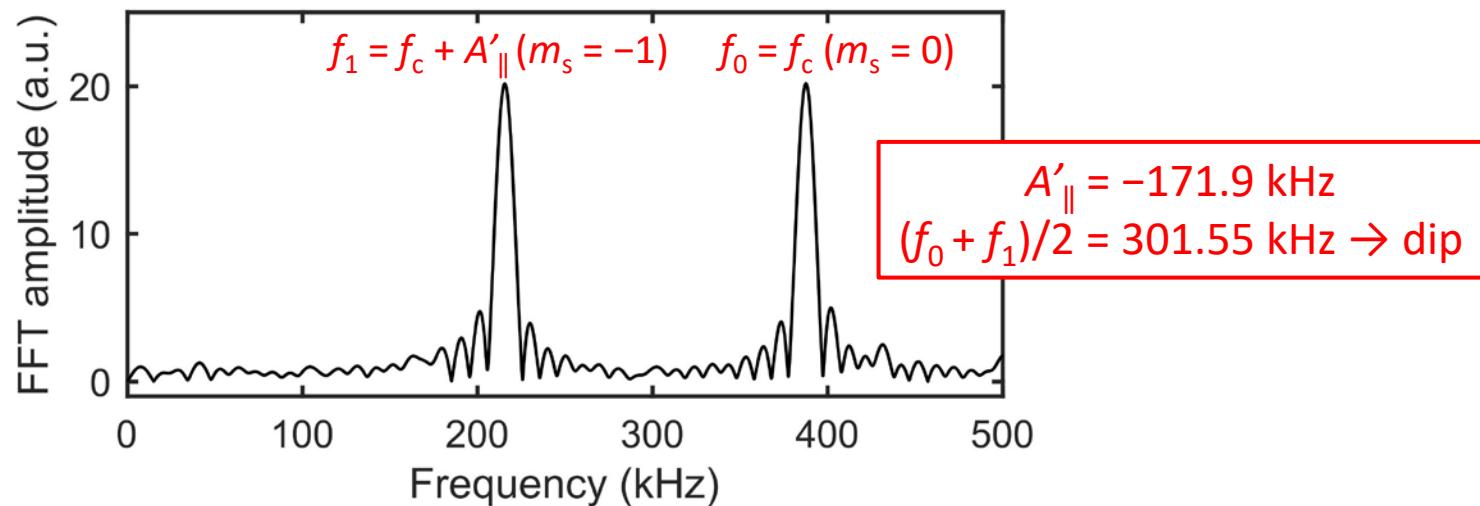
# Correlation spectroscopy of a nucleus



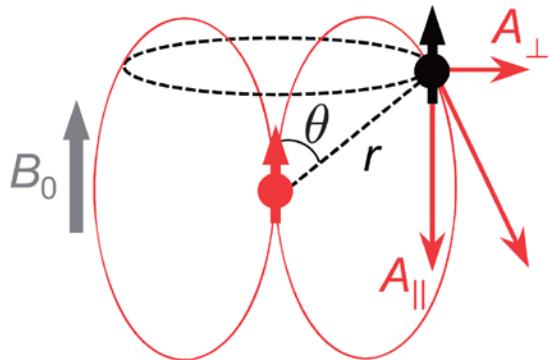
Hamiltonian of NV-<sup>13</sup>C coupled system

$$H = f_c I_z + |m_s = -1\rangle \langle -1| (A_{\parallel} I_z + A_{\perp} I_x)$$

→ No hyperfine field when  $|m_s = 0\rangle$



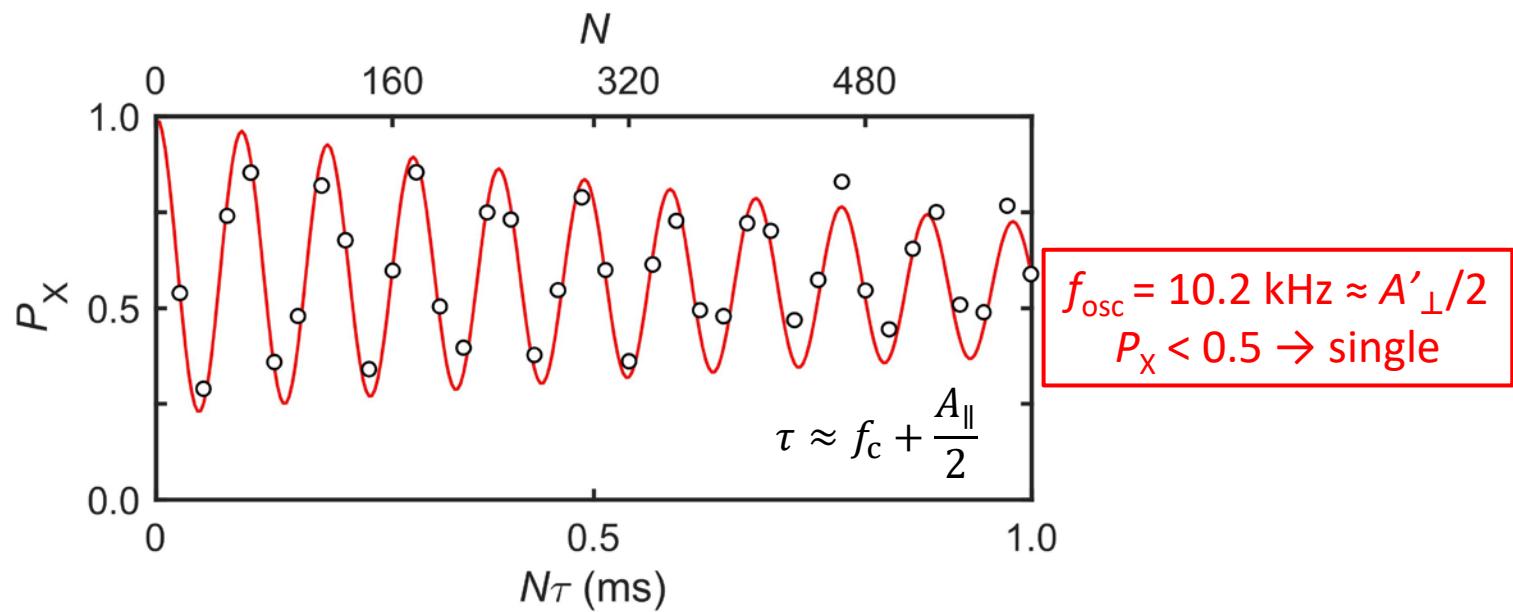
# Coherent control of a nuclear spin



Hamiltonian of NV-<sup>13</sup>C coupled system

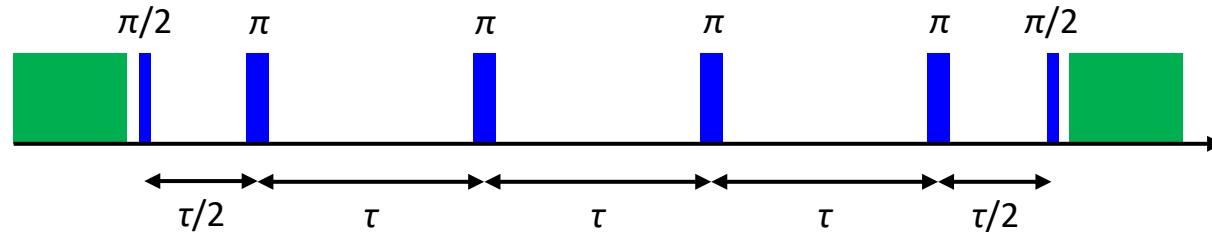
$$H = f_c I_z + |m_s = -1\rangle \langle -1| (A_{\parallel} I_z + A_{\perp} I_x)$$

→ The single <sup>13</sup>C n-spin rotates about the  $A_{\perp}$  axis

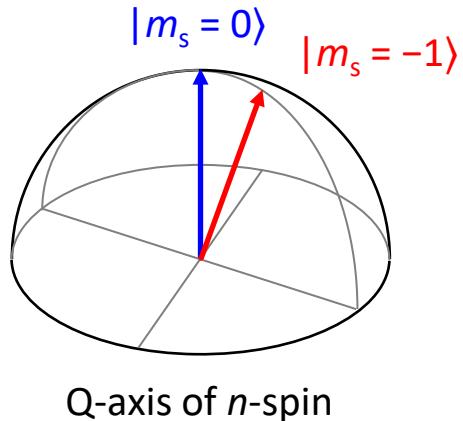
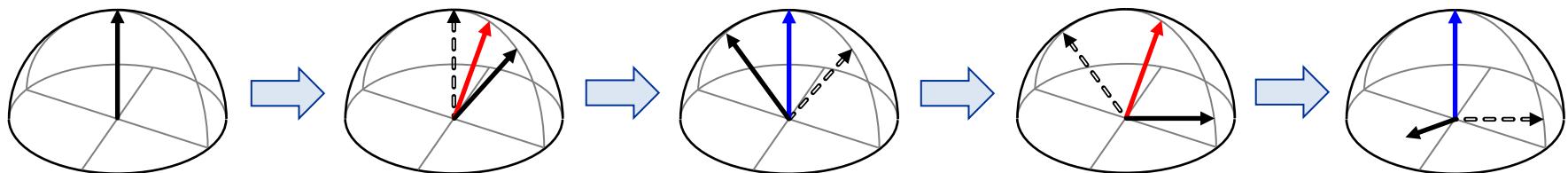


# Conditional rotation of a nuclear spin

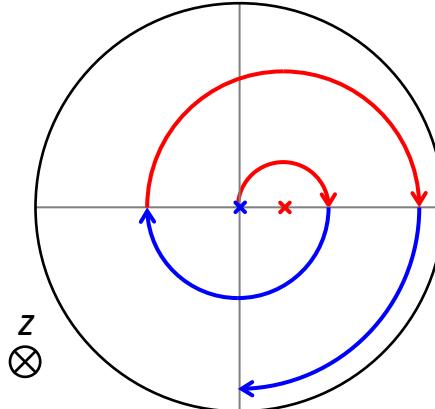
CP ( $N = 4$ )



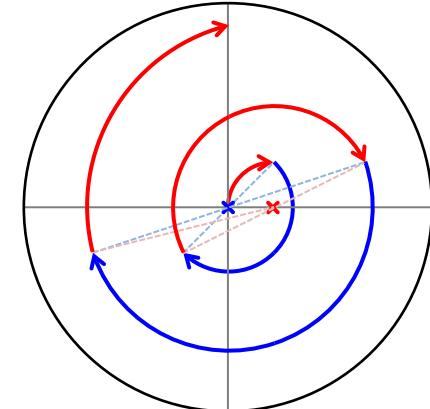
Evolution of  $n$ -spin vector



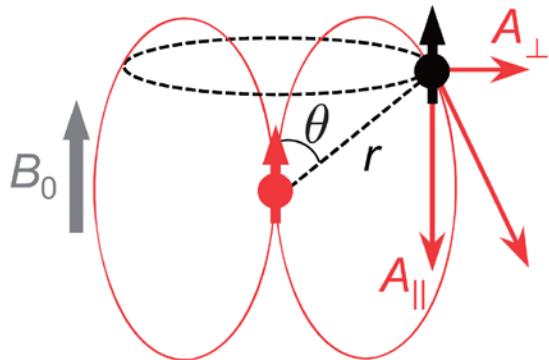
*Start from  $|m_s = 0\rangle$*



*Start from  $|m_s = -1\rangle$*



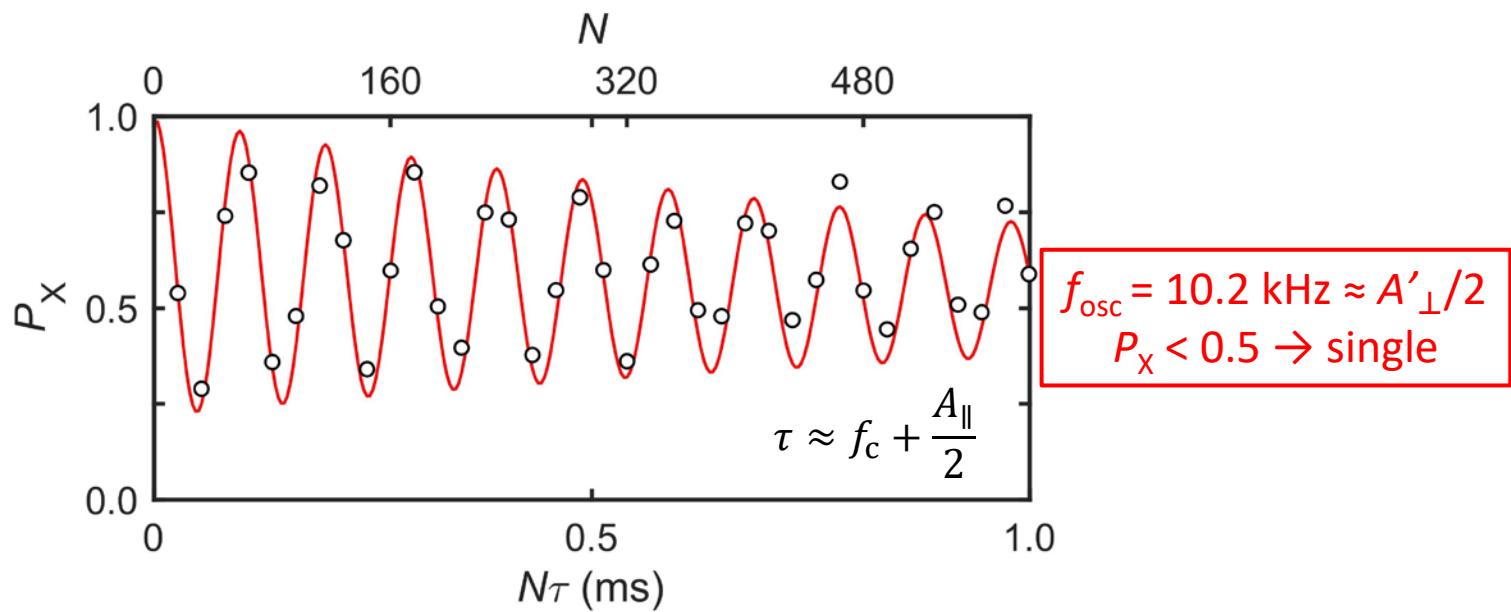
# Coherent control of a nuclear spin



Transition probability of the NV spin

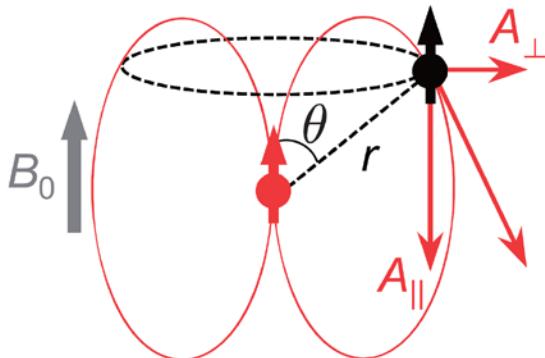
$$P_X = 1 - \frac{1}{2} \underbrace{(1 - \mathbf{n}_0 \cdot \mathbf{n}_{-1})}_{-1} \sin^2 \frac{N\phi_{cp}}{2}$$

Phys. Rev. Lett. **109**, 137602 (2012) Taminiau *et al.*



Phys. Rev. B **98**, 121405 (2018) Sasaki *et al.*

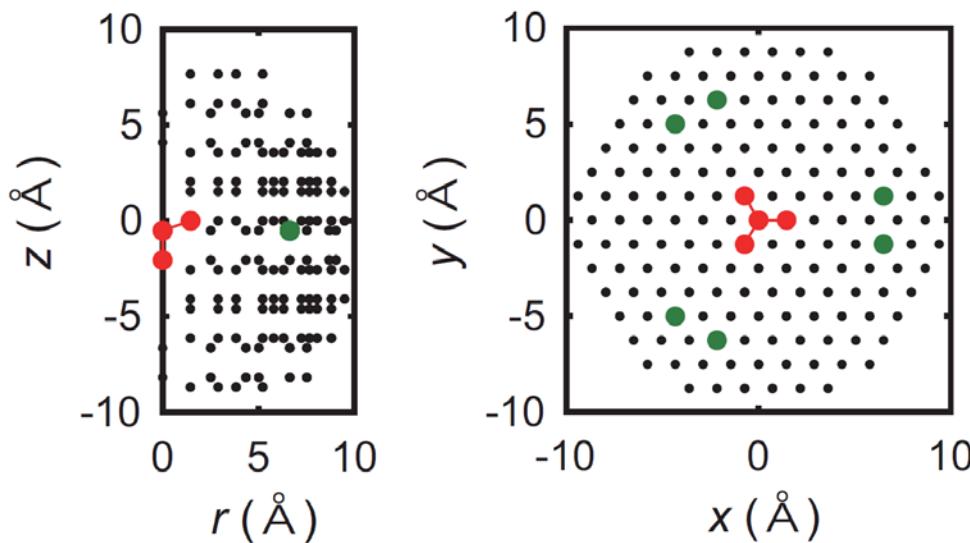
# Determination of hf constants



Magnetic dipole int. + contact hf int.

$$A_{\parallel} \propto \frac{3 \cos^2 \theta - 1}{r^3}$$

$$A_{\perp} \propto \frac{3 \cos \theta \sin \theta}{r^3}$$



$$(r, \theta) = (6.84 \text{ \AA}, 94.8^\circ)$$

$$\begin{aligned} A_{\parallel} &= -173.1 \text{ kHz} \\ A_{\perp} &= 22.3 \text{ kHz} \end{aligned}$$

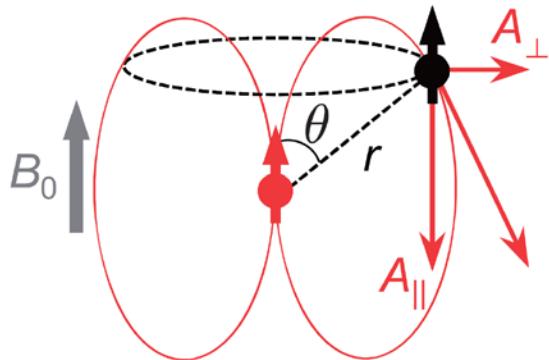
$$\begin{aligned} A_{\parallel} &= -175.1 \pm 2.1 \text{ kHz} \\ A_{\perp} &= 21.9 \pm 0.2 \text{ kHz} \end{aligned}$$

DFT: New J. Phys. **20**, 023022 (2018)  
Nizovtsev *et al.*

Phys. Rev. B **98**, 121405 (2018) Sasaki *et al.*

# How to determine $\phi$ ?

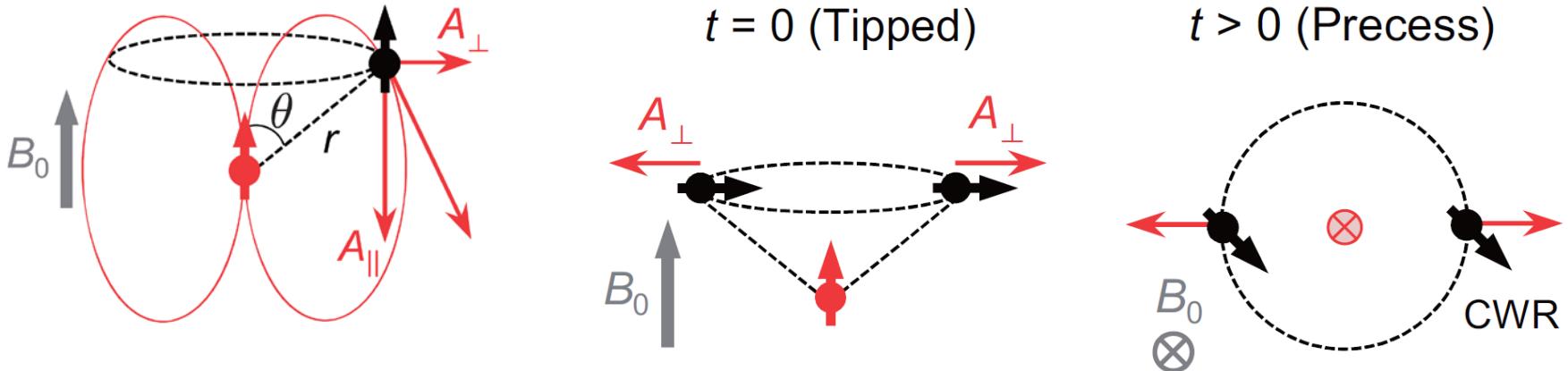
(azimuthal angle)



Magnetic dipole int.

$$A_{||} \propto \frac{3 \cos^2 \theta - 1}{r^3} \quad A_{\perp} \propto \frac{3 \cos \theta \sin \theta}{r^3}$$

# How to determine $\phi$ ?



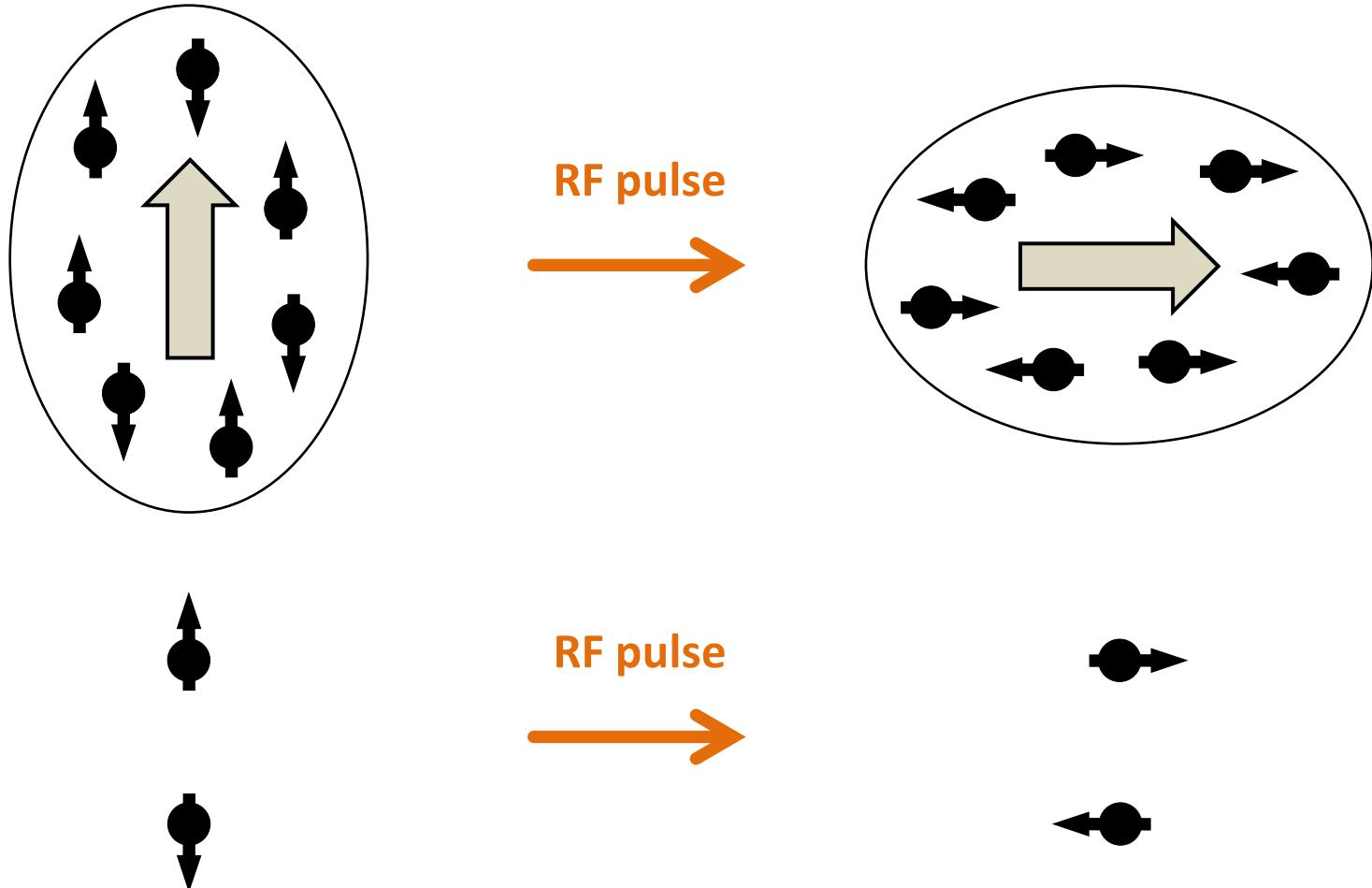
Transition probability of the NV spin after the detection of a single nuclear spin

$$P_Y = \frac{1}{2} - \frac{1}{2} \cos(\phi - \phi_n) \sin N\phi_{cp}$$



Azimuthal angle of the nuclear Bloch vector:  $2\pi f_p t + \phi_n(0)$

# Ensemble vs. single

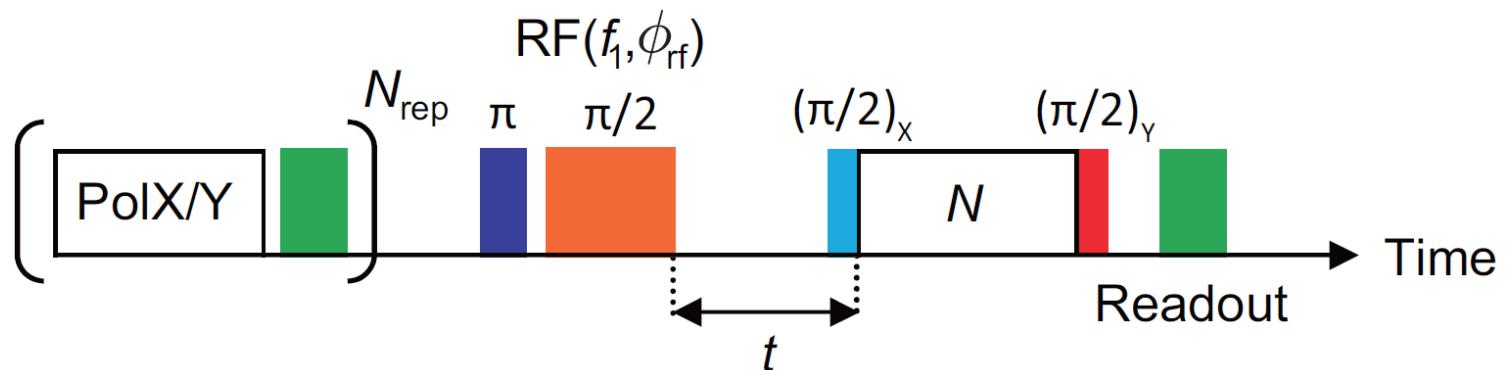
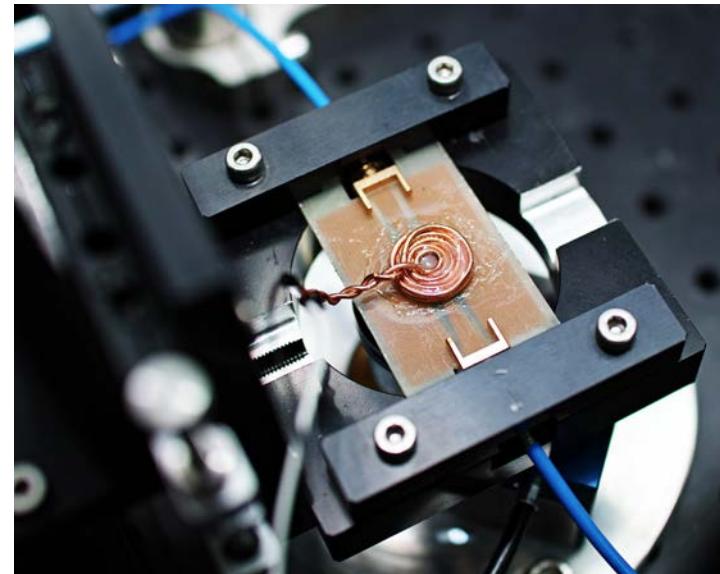


**The initial state matters**

→ Dynamic nuclear polarization (DNP)

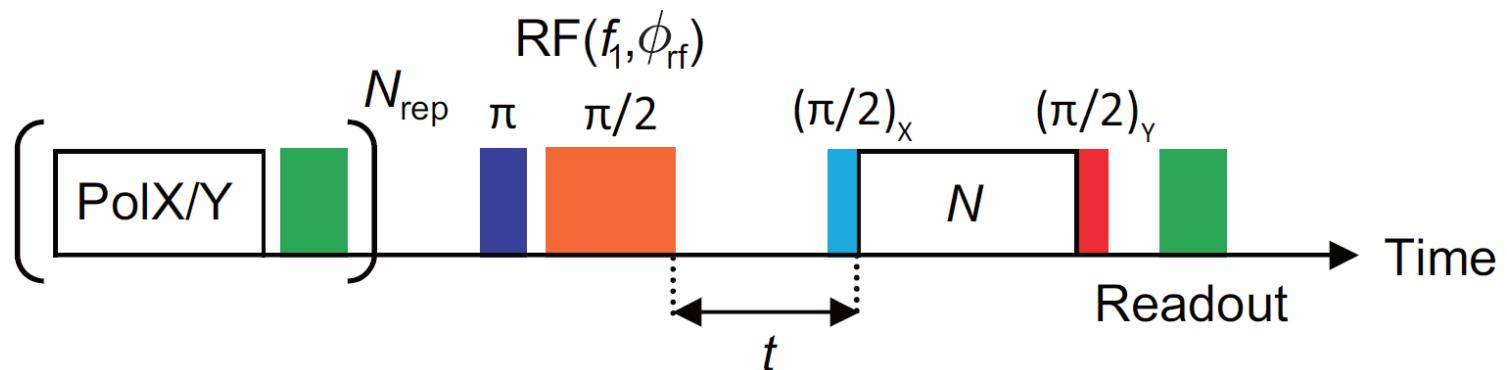
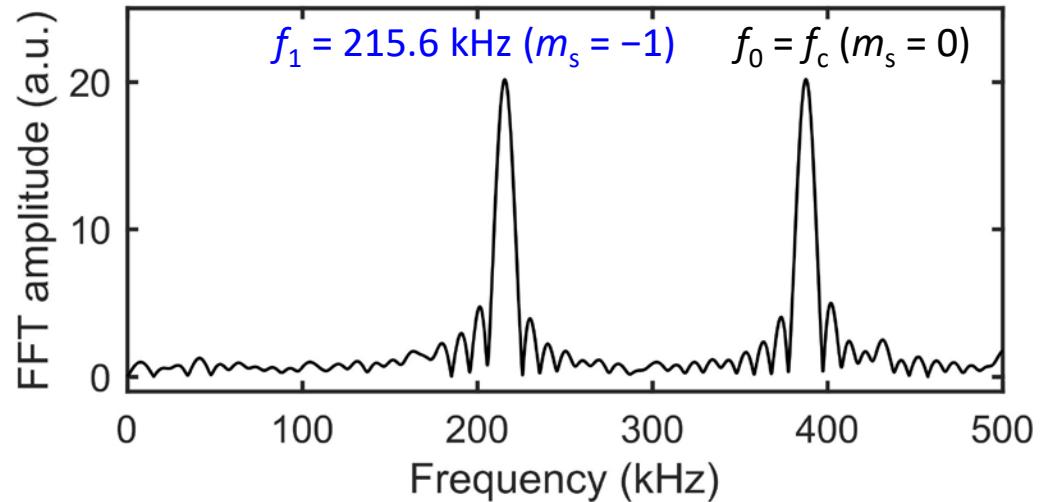
# Determination of $\phi$ of a $^{13}\text{C}$ $n$ -spin

1. DNP (PulsePol)
2. RF pulse@ $m_s = -1$
3. Wait  $t$  ( $n$ -spin precesses)
4. AC sensing



# Determination of $\phi$ of a $^{13}\text{C}$ $n$ -spin

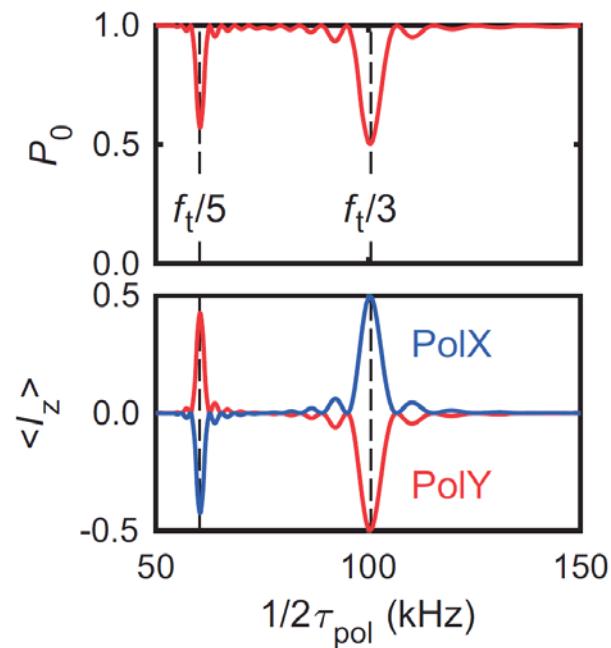
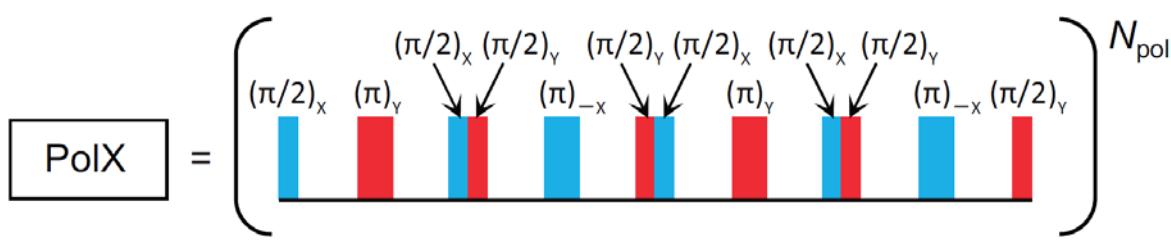
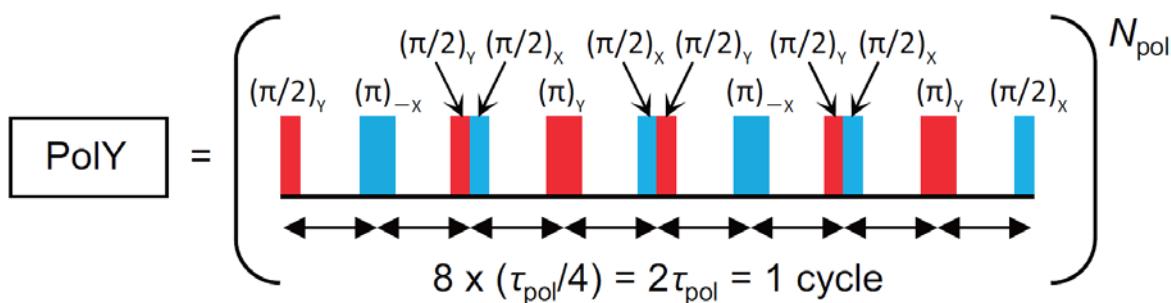
1. DNP (PulsePol)
2. RF pulse@ $m_s = -1$
3. Wait  $t$  ( $n$ -spin precesses)
4. AC sensing



# PulsePol

## Hamiltonian engineering

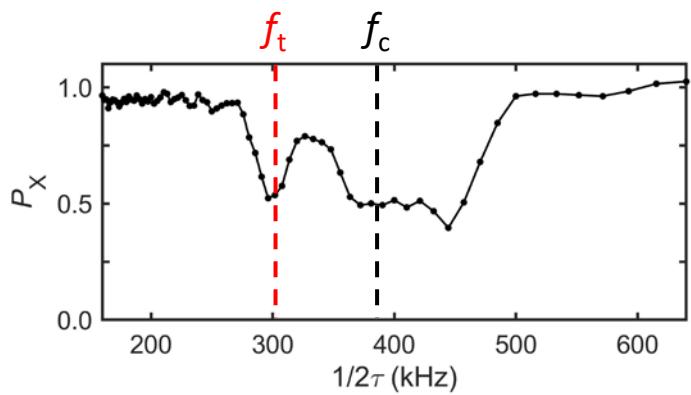
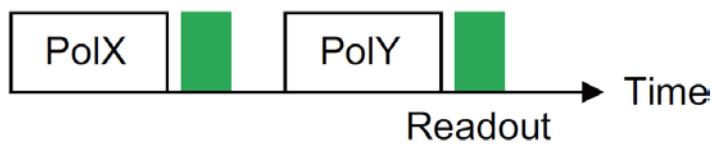
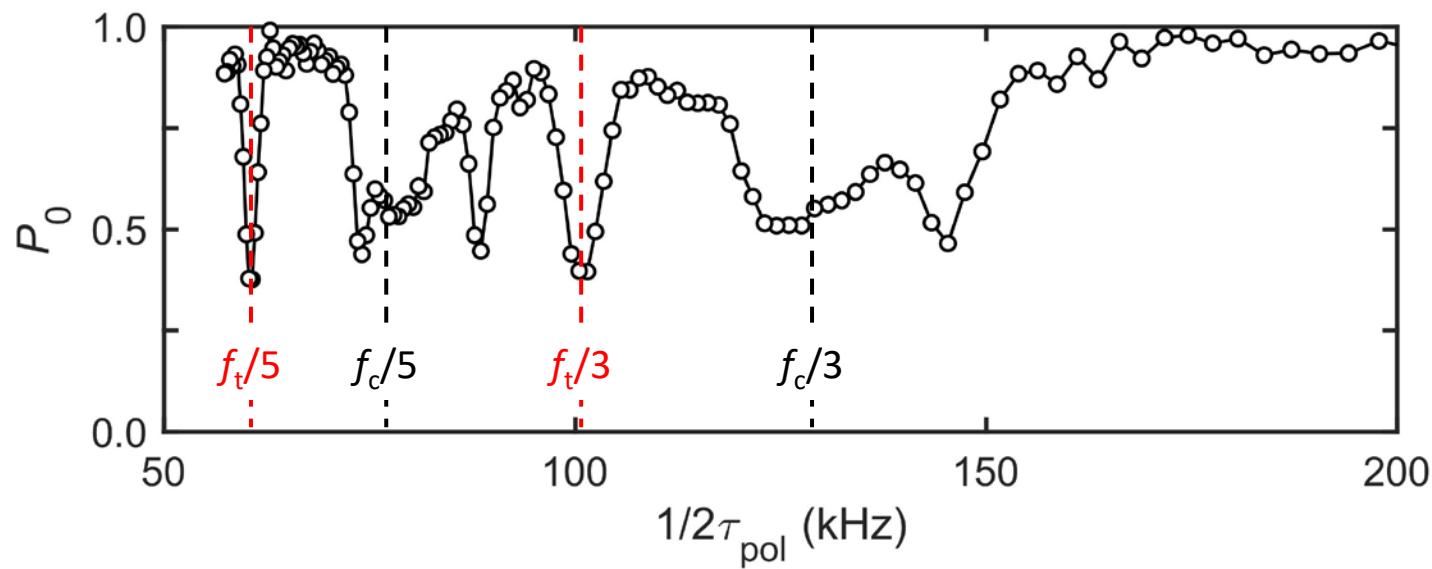
- Average Hamiltonian  $\propto S_+I_- + S_-I_+$ ,  $\propto S_+I_+ + S_-I_-$
- DNP condition:  $1/(2\tau_{\text{pol}}) = f_n/k$  ( $f_n$ :  $n$ -precession frequency,  $k$ : odd)



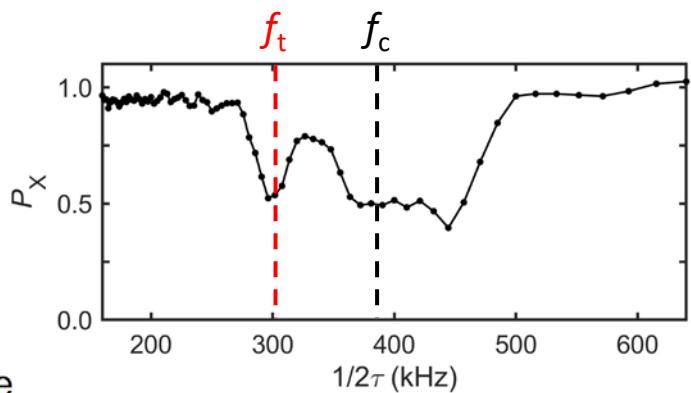
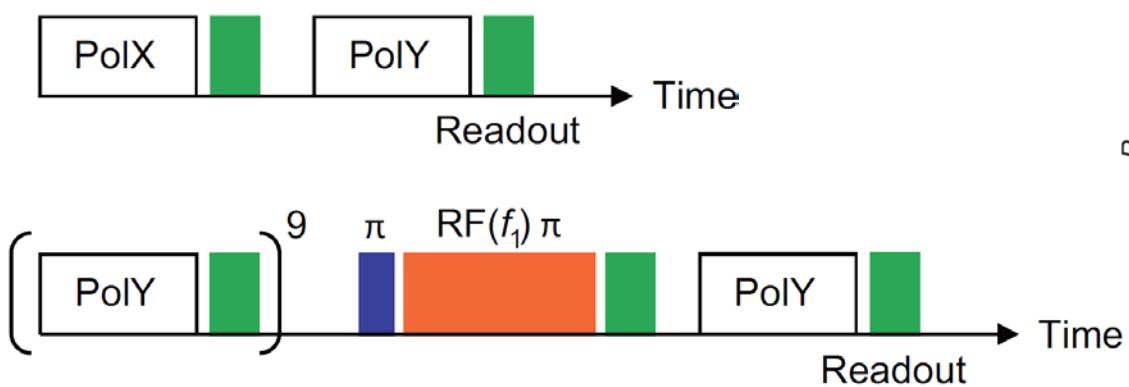
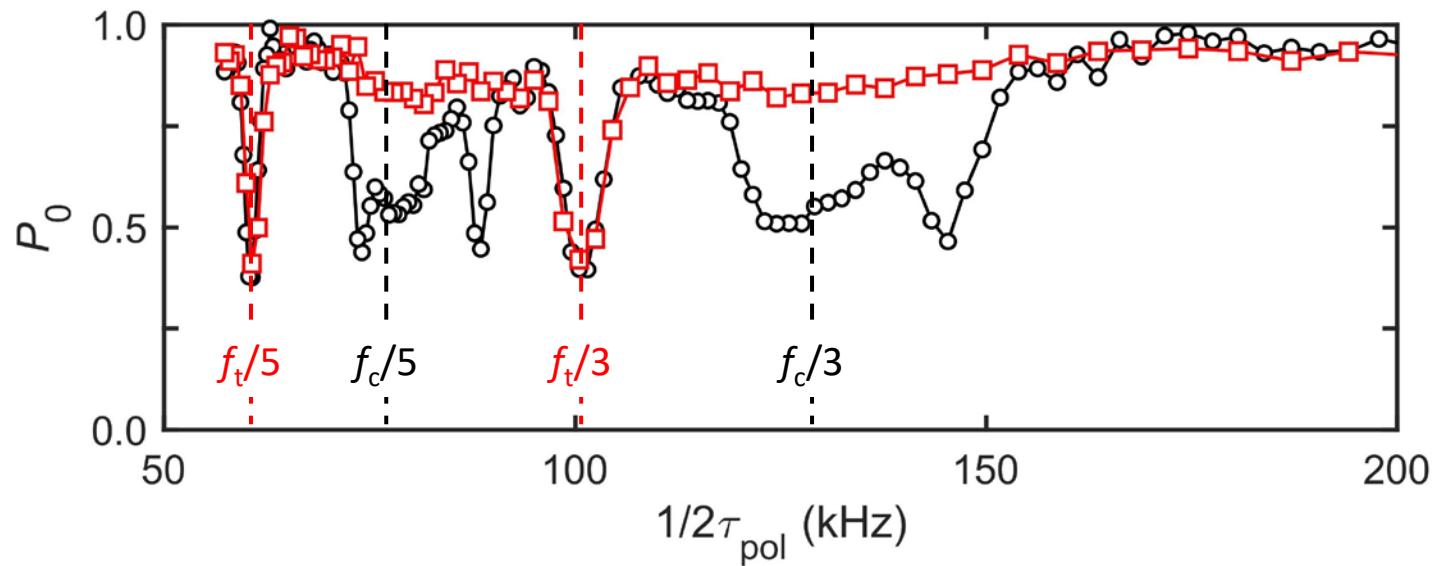
Sci. Adv. **4**, eaat8978 (2018) Schwartz *et al.*

Phys. Rev. B **98**, 121405 (2018) Sasaki *et al.*

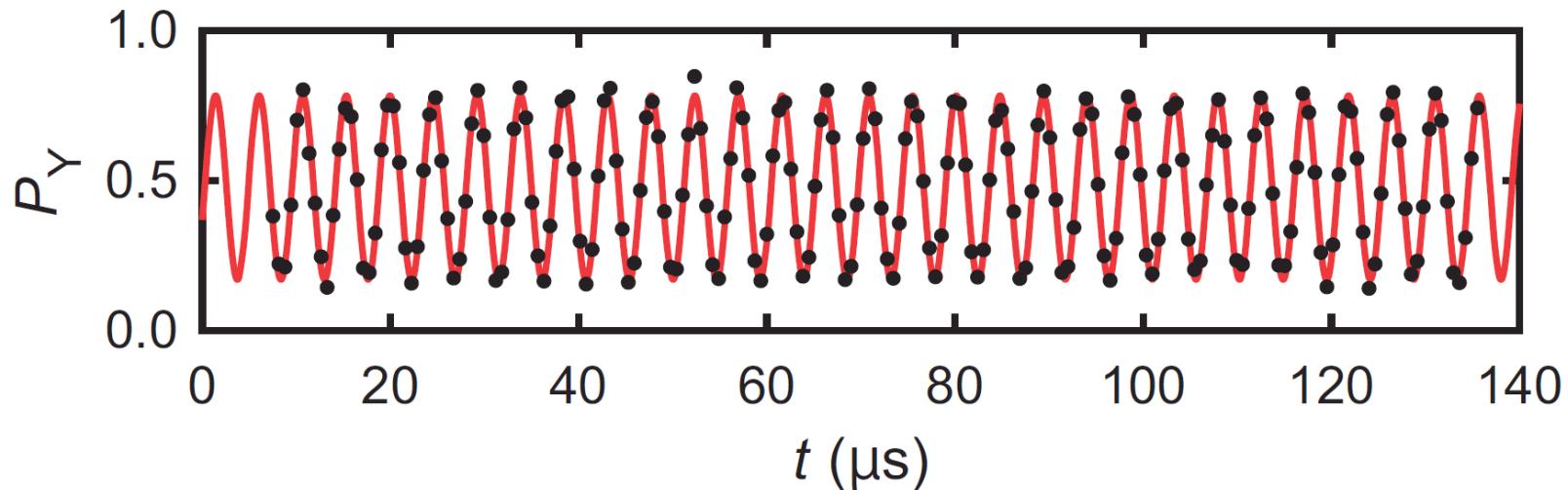
# PulsePol



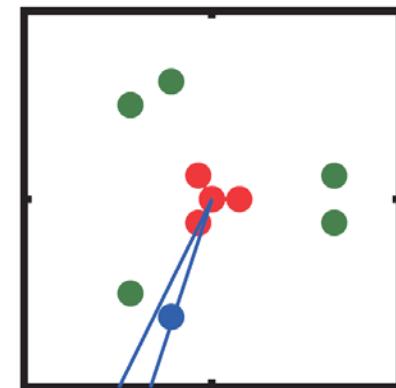
# PulsePol



# Determination of $\phi$ of a $^{13}\text{C}$ $n$ -spin

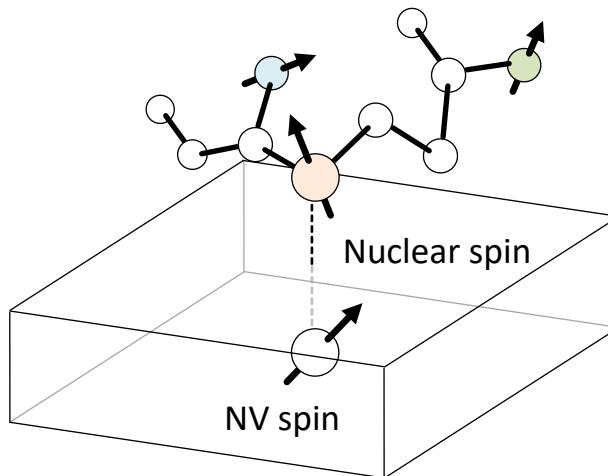


- ✓  $t \rightarrow 1 \text{ ms}$  (undersampling)
  - ✓  $f_p = 215.79 \text{ kHz} \approx f_1 = 215.6 \text{ kHz}$
  - ✓  $\phi - \phi_n(0) = 334.0^\circ$
  - ✓  $\phi_n(0) = 89.2^\circ$  (Real-space  $n$ -spin trajectory)
- $\rightarrow \phi = 247.8 \pm 4.1^\circ$



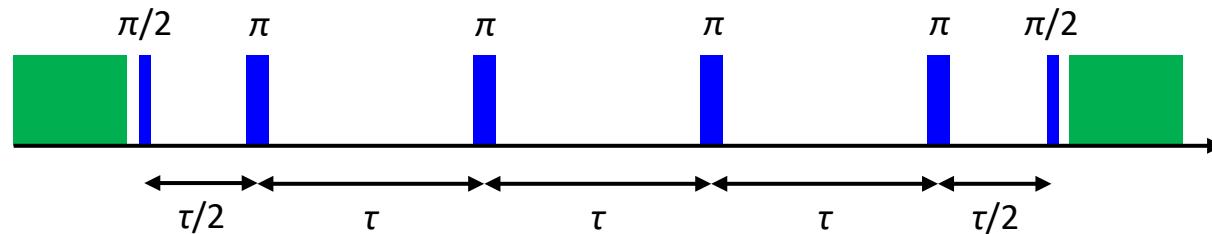
# Toward single-molecular imaging

- **Information of the positions of the individual nuclei**
  - Accurate measurement of  $e-n$  int. const's ( $A_{\parallel}, A_{\perp}$ )  $\approx (r, \theta)$
  - Lack of information on the azimuthal angle  $\phi$
- **Spectral resolution**
  - Easy to resolve isotopes
  - Need to measure  $J$ -couplings & chemical shifts (ppm!)
  - Limited by sensor/memory lifetimes ( $T_{2e/n}, T_{1e/n}$ )

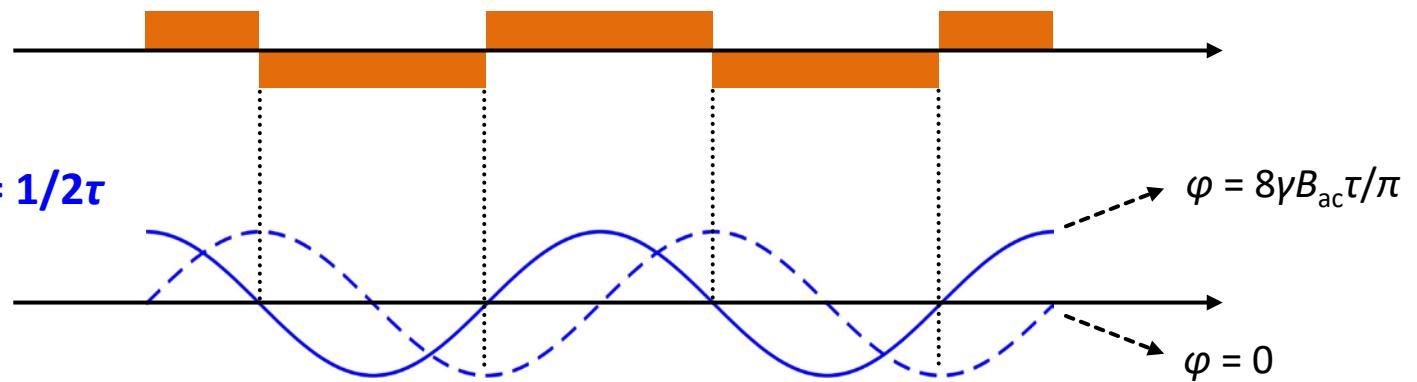


# AC magnetometry

CP ( $N = 4$ )



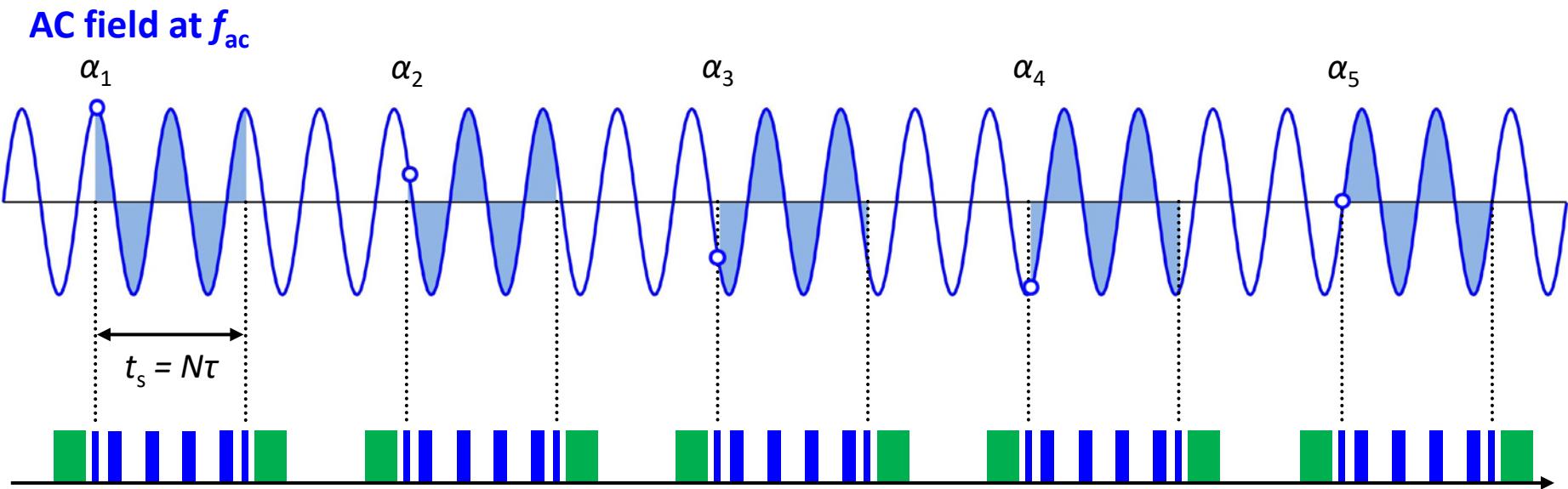
Modulation function



AC field at  $f_{ac} = 1/2\tau$

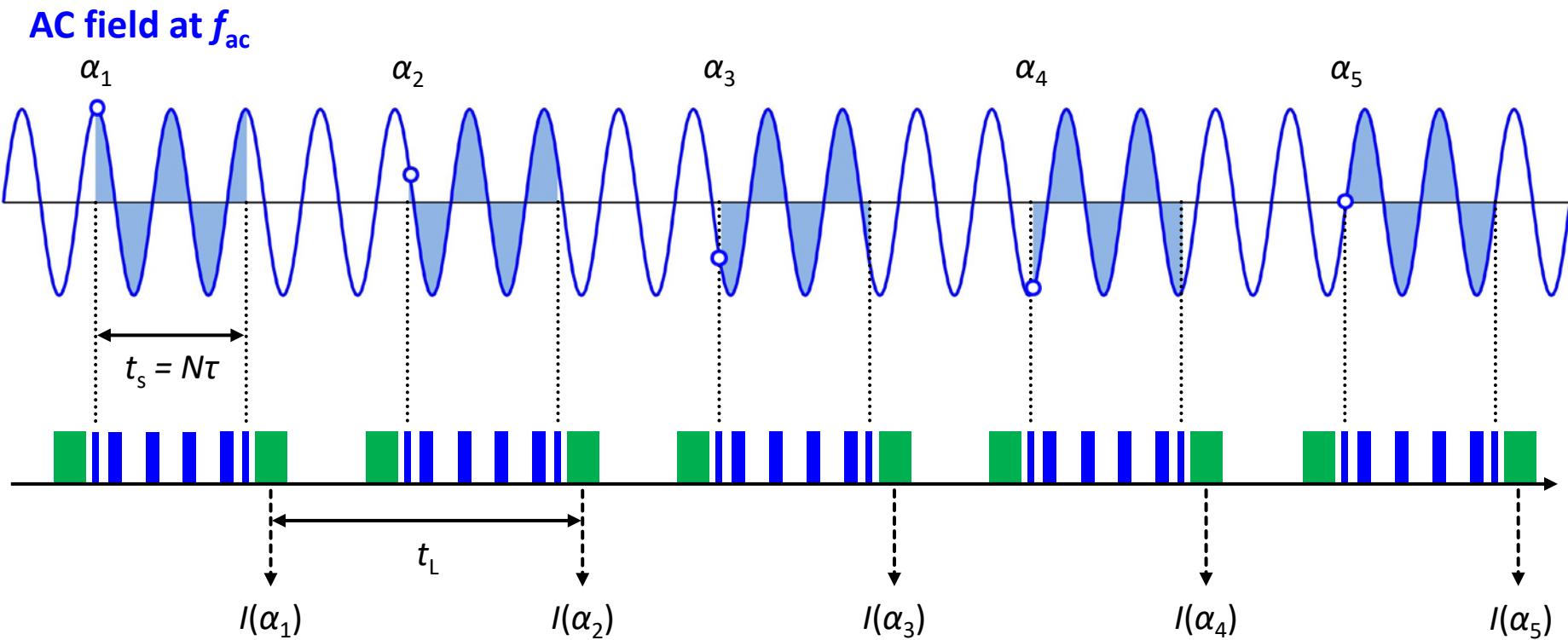
- $\varphi$  depends on the **initial phase  $\alpha$  of the AC field** ( $\varphi \propto \cos \alpha$ )

# AC magnetometry



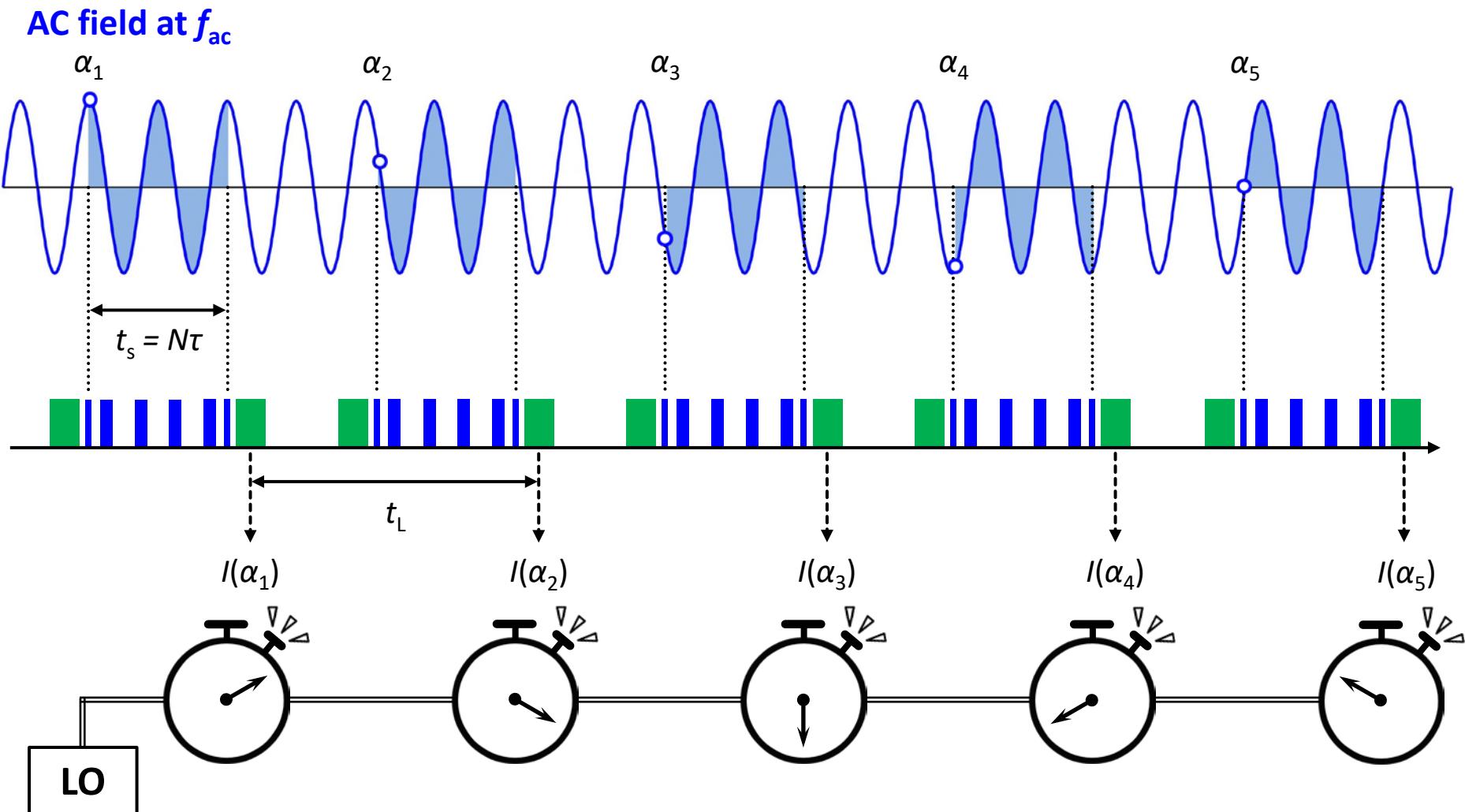
- $\varphi$  depends on the **initial phase  $\alpha$  of the AC field** ( $\varphi \propto \cos \alpha$ )
- Average over **random  $\alpha$**

# Ultrahigh resolution sensing



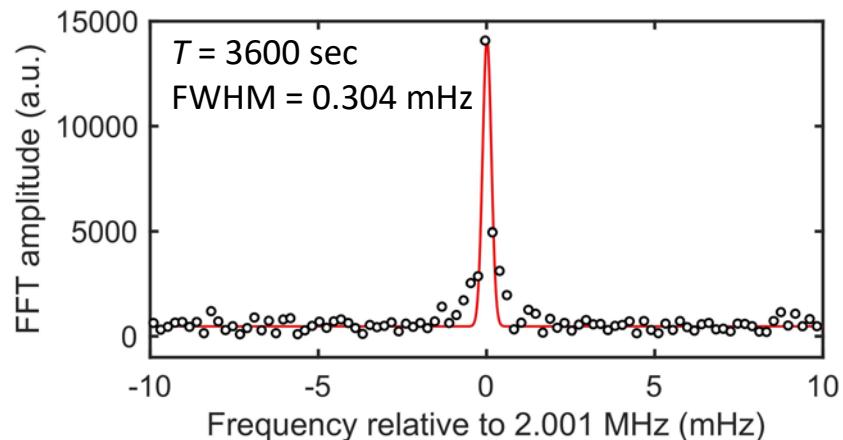
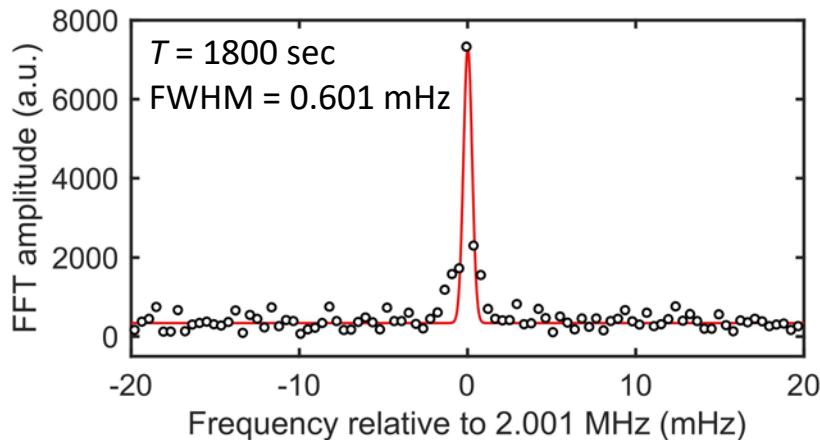
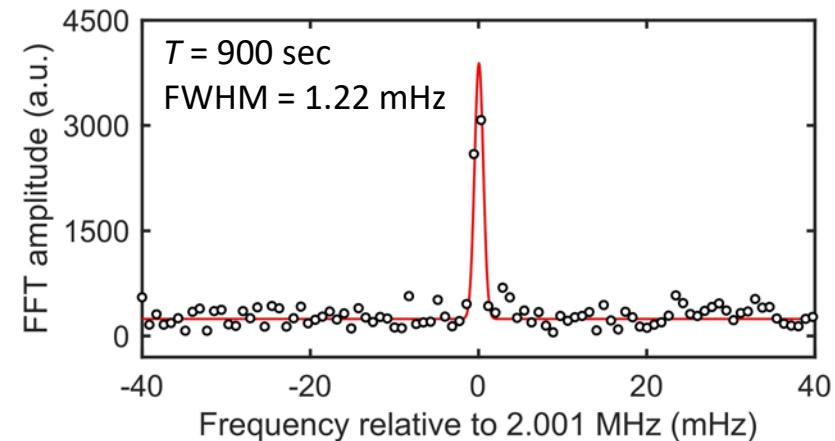
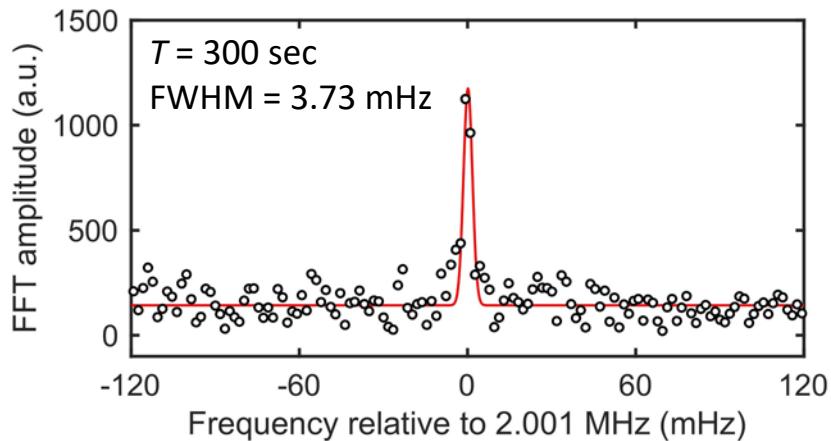
- $\varphi$  depends on the **initial phase  $\alpha$  of the AC field** ( $\varphi \propto \cos \alpha$ )
- Average over **random  $\alpha$**
- **If the data acq. is periodic**, adjacent  $\alpha$ 's are related by  $\alpha_{k+1} = 2\pi f_{ac} t_L + \alpha_k$

# Ultrahigh resolution sensing

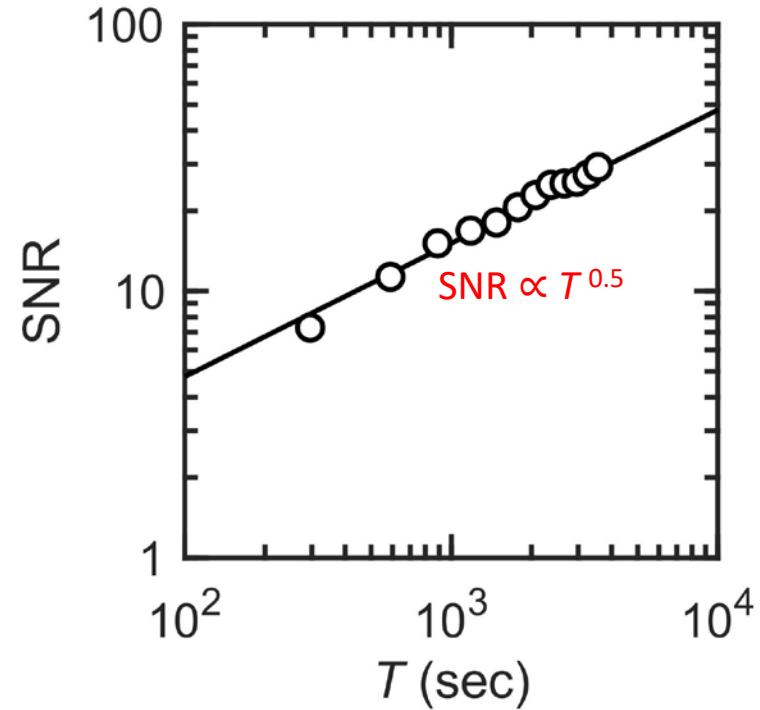
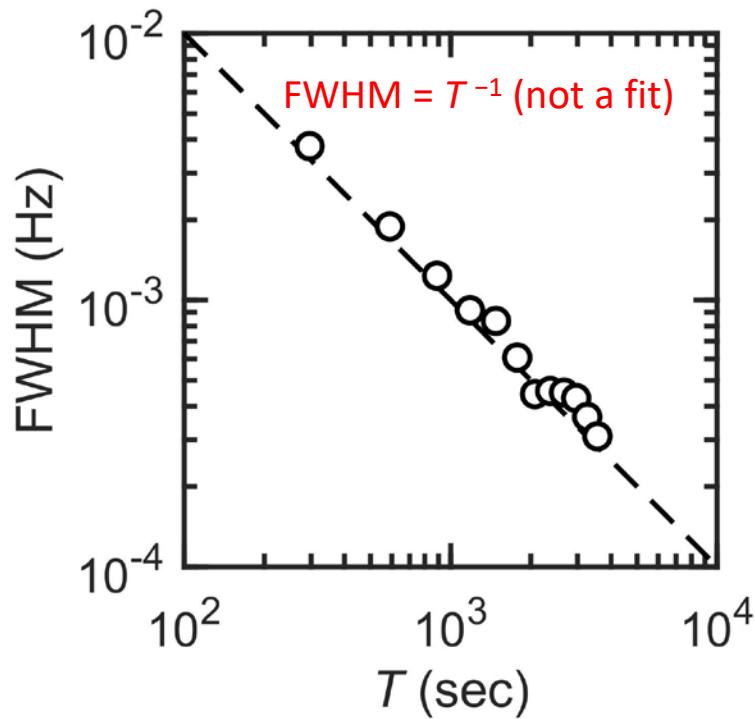


# Ultrahigh resolution sensing

$B_{\text{ac}} = 96.5 \text{ nT}$  &  $f_{\text{ac}} = 2.001 \text{ MHz}$  applied from a coil, detected by a single NV center



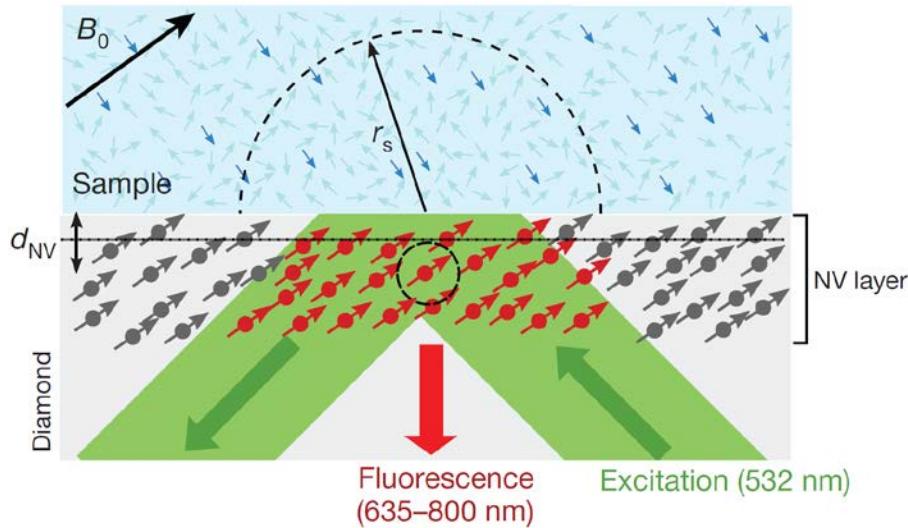
# Ultrahigh resolution sensing



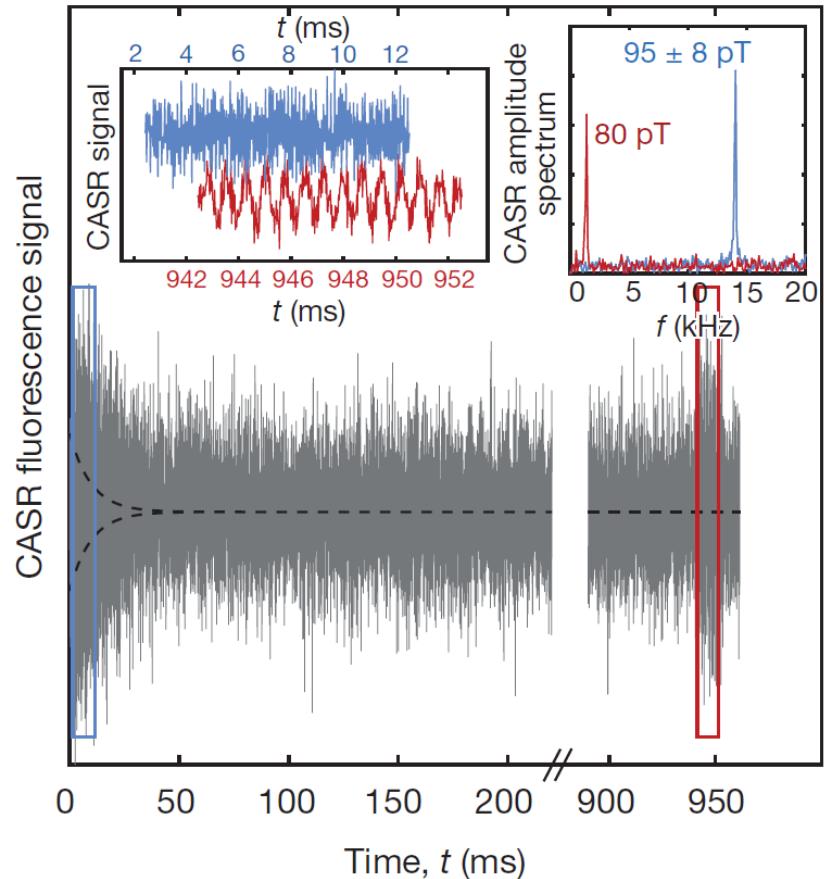
- Spectral resolution not limited by sensor/memory lifetimes ( $T_{2e/n}$ ,  $T_{1e/n}$ )
- Only limited by the stability of LO (essentially infinite)
- Resolution =  $T^{-1}$  & SNR  $\propto T^{0.5} \rightarrow$  Precision  $\propto T^{-1.5}$

# NMR spectroscopy

Data from Harvard: Nature **555**, 351 (2018) Glenn *et al.*



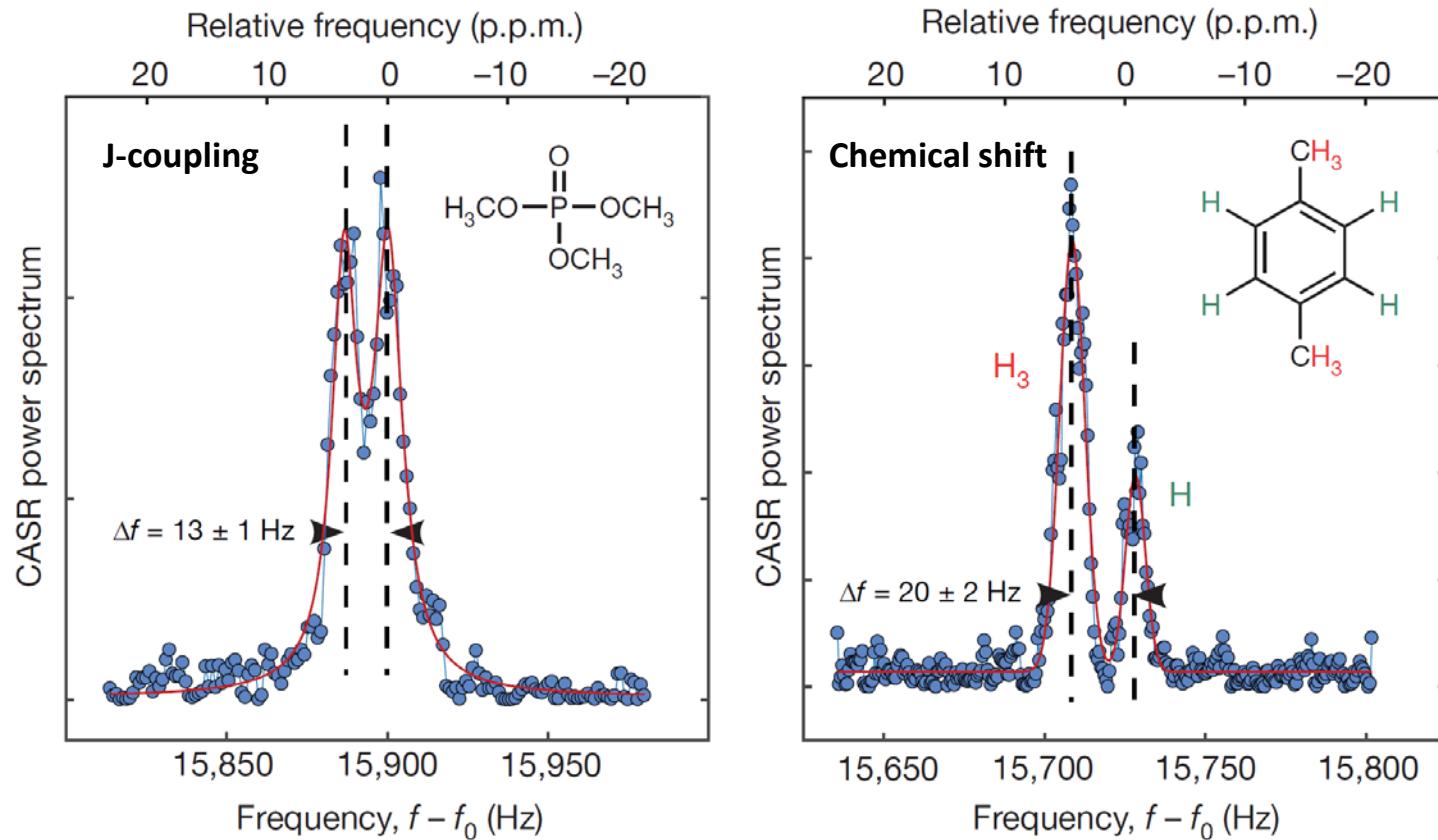
- $[NV] \approx 3 \times 10^{17} \text{ cm}^{-3}$
- # of NV  $\approx 5 \times 10^9$
- $V_{\text{detect}} \approx 25 \text{ pL}$
- # of protons  $\approx 2.5 \times 10^{15}$
- RF pulse  $\rightarrow$  FID



See also: Science **357**, 67 (2017) Aslam *et al.* (Wrachtrup, Stuttgart)  
[ $B_0 = 3 \text{ T}$ ,  $f_e = 87 \text{ GHz}$ ,  $T_{1n} = 260 \text{ s}$ ]

# NMR spectroscopy

Data from Harvard: Nature **555**, 351 (2018) Glenn *et al.*



See also: Science **357**, 67 (2017) Aslam *et al.* (Wrachtrup, Stuttgart)  
[ $B_0 = 3$  T,  $f_e = 87$  GHz,  $T_{1n} = 260$  s]

# Summary

- **Tools for single-molecule imaging/structural analysis are being developed**
  - Determination of the position of individual  $n$ -spins<sup>[1,2,3]</sup>
  - Ultrahigh resolution sensing<sup>[4,5,6]</sup>, resolving chemical shifts<sup>[6,7]</sup> & suppression of backaction from  $n$ -spins<sup>[8,9]</sup>

[1] Phys. Rev. B **98**, 121405 (2018) Sasaki *et al.* (Keio)

[2] Phys. Rev. Lett. **121**, 170801 (2018) Zopes *et al.* (ETH)

[3] Nature **576**, 411 (2019) Abobeih *et al.* (Delft)

[4] Science **356**, 832 (2017) Schmitt *et al.* (Ulm)

[5] Science **356**, 837 (2017) Boss *et al.* (ETH)

[6] Nature **555**, 351 (2018) Glenn *et al.* (Harvard)

[7] Science **357**, 67 (2017) Aslam *et al.* (Stuttgart)

[8] Nature Commun. **10**, 594 (2019) Pfender *et al.* (Stuttgart)

[9] Nature **571**, 230 (2019) Cujia *et al.* (ETH)